

# The *Suzaku* view of highly ionized outflows in AGN – I. Statistical detection and global absorber properties

Jason Gofford,<sup>1</sup>★ James N. Reeves,<sup>1,2</sup> Francesco Tombesi,<sup>3,4</sup> Valentina Braito,<sup>5,6</sup>  
T. Jane Turner,<sup>2</sup> Lance Miller<sup>7</sup> and Massimo Cappi<sup>8</sup>

<sup>1</sup>*Astrophysics Group, Keele University, Keele ST5 5BG, UK*

<sup>2</sup>*Department of Physics, University of Maryland Baltimore County, Baltimore, MD 21250, USA*

<sup>3</sup>*X-ray Astrophysics Laboratory and CRESST, NASA/GSFC, Greenbelt, MD 20771, USA*

<sup>4</sup>*Department of Astronomy, University of Maryland, College Park, MD 20742, USA*

<sup>5</sup>*INAF-Osservatorio Astronomico di Brera, Via Bianchi 46, I-23807 Merate, Italy*

<sup>6</sup>*Department of Physics and Astronomy, University of Leicester, University Road, Leicester LE1 7RH, UK*

<sup>7</sup>*Department of Physics, Oxford University, Denys Wilkinson Building, Keble Road, Oxford OX1 3RH, UK*

<sup>8</sup>*INAF-IASF Bologna, Via Gobetti 101, I-40129 Bologna, Italy*

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## ABSTRACT

We present the results of a new spectroscopic study of Fe K-band absorption in active galactic nuclei (AGN). Using data obtained from the *Suzaku* public archive we have performed a statistically driven blind search for Fe xxv He $\alpha$  and/or Fe xxvi Ly $\alpha$  absorption lines in a large sample of 51 Type 1.0–1.9 AGN. Through extensive Monte Carlo simulations we find that statistically significant absorption is detected at  $E \gtrsim 6.7$  keV in 20/51 sources at the  $P_{\text{MC}} \geq 95$  per cent level, which corresponds to  $\sim 40$  per cent of the total sample. In all cases, individual absorption lines are detected independently and simultaneously amongst the two (or three) available X-ray imaging spectrometer detectors, which confirms the robustness of the line detections. The most frequently observed outflow phenomenology consists of two discrete absorption troughs corresponding to Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  at a common velocity shift. From *xstar* fitting the mean column density and ionization parameter for the Fe K absorption components are  $\log(N_{\text{H}}/\text{cm}^{-2}) \approx 23$  and  $\log(\xi/\text{erg cm s}^{-1}) \approx 4.5$ , respectively. Measured outflow velocities span a continuous range from  $< 1500$  km s $^{-1}$  up to  $\sim 100\,000$  km s $^{-1}$ , with mean and median values of  $\sim 0.1$  c and  $\sim 0.056$  c, respectively. The results of this work are consistent with those recently obtained using *XMM-Newton* and independently provides strong evidence for the existence of very highly ionized circumnuclear material in a significant fraction of both radio-quiet and radio-loud AGN in the local universe.

**Key words:** line: identification – galaxies: active – galaxies: nuclei – X-rays: galaxies.

## 1 INTRODUCTION

Observational evidence for outflows and winds in active galactic nuclei (AGN) is seen in multiple energy regimes, ranging from the prominent radio-jets seen in radio-loud sources, to the broad absorption lines observed in broad absorption line (BAL) quasars, through to the photoionized ‘warm absorber’ which is frequently observed in the soft X-rays (e.g. Crenshaw, Kraemer & George 2003; Blustin et al. 2005; McKernan, Yaqoob & Reynolds 2007). Gravitational micro-lensing studies have shown that the primary X-ray emission region in AGN is of the order of a few tens of

gravitational radii ( $R_g = GM/c^2$ ) in size (Morgan et al. 2008; Chartas et al. 2009; Dai et al. 2010) and so the spectral analysis of absorption features imprinted on the X-ray continuum by circumnuclear material in an AGN is a powerful diagnostic of the physical conditions of the environment in the vicinity of the central super-massive black hole (SMBH), of the dynamics and kinematics of the outflowing material, its chemical composition and its ionization state. Understanding how such winds are formed, their physical characteristics and their energetic output are of vital importance when it comes to studying how the interplay between the accretion and ejection flows at small radii can affect the host galaxy on larger scales (e.g. Ferrarese & Merritt 2000; King & Pounds 2003; Sazonov et al. 2005; Ciotti, Ostriker & Proga 2009).

★ E-mail: jag@astro.keele.ac.uk

In the X-ray regime the presence of photoionized material in AGN is well established, with at least 50 per cent of objects showing direct spectroscopic evidence for discrete absorption lines and photoelectric edges in their soft-band ( $E < 3$  keV) spectra (e.g. Reynolds 1997; Crenshaw et al. 2003; Blustin et al. 2005), with typical line-of-sight velocities ranging from a few hundred to around a thousand  $\text{km s}^{-1}$  (Blustin et al. 2005; McKernan et al. 2007). Detailed modelling at high spectral resolution with photoionization codes such as *xstar* (Kallman et al. 2004) often finds the soft-band absorber to be described by column densities and ionization parameters<sup>1</sup> in the range of  $\log(N_{\text{H}}/\text{cm}^{-2}) \sim 20\text{--}23$  and  $\log(\xi/\text{erg cm s}^{-1}) \sim 0\text{--}3$ , respectively. These parameters imply that the warm absorbers are typically located on fairly large distances from the central black hole, and perhaps associated with a wind originating from the putative parsec scale torus (Blustin et al. 2005) or with the latter stages of an accretion disc wind which has propagated out to larger radii (Proga & Kallman 2004). By virtue of their low outflow velocities the soft X-ray warm absorbers are thought to only have a weak feedback effect in their host galaxy. Indeed, the mechanical power imparted by individual warm absorption components is very low, typically  $\lesssim 0.01$  per cent of an AGN’s bolometric luminosity ( $L_{\text{bol}}$ ) (e.g. Blustin et al. 2005), which is significantly lower than the  $\sim 0.5$  per cent of  $L_{\text{bol}}$  thought necessary for feedback to affect the host galaxy (Hopkins & Elvis 2010). However, Crenshaw & Kraemer (2012) have recently shown that this  $\sim 0.5$  per cent threshold can be exceeded provided the mechanical power is integrated over all UV and X-ray absorption components, at least in the case of a few moderate-luminosity local AGN.

More recently the higher throughput and larger effective area offered by *XMM-Newton* and *Suzaku* at higher X-ray energies ( $E \sim 5\text{--}10$  keV) have shown that absorption, specifically in the form of very highly ionized resonant absorption lines associated with the K-shell ( $1s \rightarrow 2p$ ) transitions of Fe  $\text{xxv}$  and Fe  $\text{xxvi}$ , is also manifested in the hard X-ray spectrum of a significant fraction of local AGN. While evidence for such absorption lines was initially confined to detailed studies of individual sources (e.g. Pounds et al. 2003b; Reeves, O’Brien & Ward 2003; Risaliti et al. 2005; Turner et al. 2008; Cappi et al. 2009; Giustini et al. 2011) the recent systematic archival *XMM-Newton* study conducted by Tombesi et al. (2010a, 2011a, 2012) has shown that Fe  $\text{xxv}\text{--}\text{xxvi}$  absorption lines are present in the X-ray spectra of  $\gtrsim 40$  per cent of radio-quiet AGN in the local universe ( $z < 0.1$ ). Moreover, such outflows have also been detected in a small sample of local broad line radio galaxies (BLRGs; Tombesi et al. 2010b, hereafter T10B) which thus suggests that they may represent an important addition to the commonly held AGN unification model (e.g. Antonucci 1993; Urry & Padovani 1995).

In comparison to the soft-band absorbers these hard X-ray absorbers generally have much more extreme parameters, with  $\log(N_{\text{H}}/\text{cm}^{-2}) \approx 23\text{--}24$  and  $\log(\xi/\text{erg cm s}^{-1}) \approx 3\text{--}6$ , and their outflow velocities relative to the host galaxy can reach mildly relativistic values. While alternative explanations have been posited in the literature, e.g. resonant absorption by highly ionized material in a corotating and optically thin plasma above the accretion disc (Gallo & Fabian 2011), the large inferred velocities – combined with the short time-scale variability sometimes exhibited by the ab-

sorption features – point to an origin more likely associated with a wind which is launched from the surface of the accretion disc itself (e.g. Pounds et al. 2003b; Reeves et al. 2009; Gofford et al. 2011; Tombesi et al. 2012). In this scenario the inferred mass outflow rates for disc winds are often comparable to those of the matter which accretes on to the central black hole and the consequent mechanical power can also be a sizeable fraction (i.e.  $\geq$  few per cent) of an AGN’s bolometric luminosity (e.g. Chartas et al. 2002; Pounds et al. 2003b; Gibson et al. 2005; Reeves et al. 2009; Gofford et al. 2011; Tombesi et al. 2012) making them a promising means of linking the small- and large-scale processes at play in galactic evolution over cosmic time.

### 1.1 Why *Suzaku*?

In this work we use archival *Suzaku* observations of a large sample of AGN to characterize the properties of highly ionized outflows in the Fe K band. To date, the most comprehensive study of these outflows has been conducted by Tombesi et al. (2010a, hereafter T10A) who performed a systematic narrow-band (i.e. 3.5–10.5 keV) analysis of 42 sources (with 101 observations) obtained from the *XMM-Newton* archive, using a simple baseline model to describe the AGN continuum in the Fe K band. The baseline model consisted of a power law, narrow Gaussians and, where required, neutral absorption to approximate for any spectral curvature. This phenomenological approach yielded a fit to the 4–10 keV energy band of most sources which was sufficient to enable the systematic search of Fe  $\text{xxv}$  He $\alpha$  and Fe  $\text{xxvi}$  Ly $\alpha$  absorption lines without needing to take additional spectral complexity into account. Moreover, T10A found that this approach resulted in continuum parameters which were largely consistent with those found by authors who conducted a more thorough fit using the entire *XMM-Newton* bandpass. However, while this approach is suitable for most sources it is important to note that in those which have more complex X-ray spectra, e.g. those which are very heavily absorbed or those with strong hard excesses, using a narrow-band fit can yield a model which is a poor representation of the data when extrapolated to consider the higher and lower energies. The only way to overcome this limitation is by performing a detailed broad-band analysis. *Suzaku* is currently the only X-ray observatory which offers a sufficiently broad bandpass (i.e. 0.6–50 keV) to allow for the effects of soft-band absorption, the soft-excess and the Compton reflection component to be constrained simultaneously. This makes it the ideal instrument to confidently constrain the underlying continuum and to robustly assess for the presence of highly ionized outflows in AGN.

## 2 SAMPLE SELECTION

The sample was selected from the Data Archive and Transmission System<sup>2</sup> (DARTS; Tamura et al. 2004) which contains all publicly available *Suzaku* observations categorized by the classification of the target source. An initial sample of AGN was drawn by positionally cross-correlating the targeted pointing coordinates of all publicly available (as of the end of 2011 December) *Suzaku* observations of extragalactic compact sources against the known positions of sources contained in the VERONCAT catalogue of Quasars & AGN (Véron-Cetty & Véron 2010). The VERONCAT has an extensive list of local AGN, all of which have been quantitatively classified based upon their optical properties using the

<sup>1</sup> The ionization parameter is defined as  $\xi = L_{\text{ion}}/nR^2$  (Tarter, Tucker & Salpeter 1969), where  $L_{\text{ion}}$  is the 1–1000 Rydberg ionizing luminosity,  $n$  is the electron density and  $R$  is the distance of the ionizing source from the absorbing clouds.

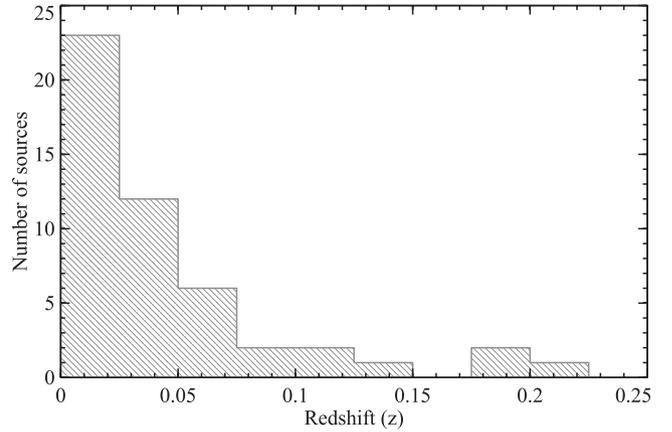
<sup>2</sup> <http://darts.jaxa.jp/astro/suzaku/>

criteria introduced by Winkler (1992). Observations of non-AGN which were matched by virtue of their having a similar position to a known AGN on the sky, such as those of extragalactic ultra-luminous X-ray sources (ULXs; Feng & Soria 2011) or X-ray bright supernovae, were systematically excluded. Only those observations with exposures long enough to ensure that the net source counts between 2 and 10 keV in the source rest-frame exceeded  $\sim 15\,000$  were retained (typical exposures were  $\gtrsim 50$  ks) such that a narrow (i.e.  $EW = 30$  eV) absorption feature was detectable at 95 per cent from Monte Carlo simulations at source rest-frame energies of up to 8–9 keV (see Section 4.4). In the case of the high red-shift quasar APM 08279+5255, which shows evidence for absorption lines at rest-frame energies of  $E > 10$  keV (e.g. Chartas et al. 2009; Saez, Chartas & Brandt 2009; Saez & Chartas 2011), the measure of net source counts was instead taken for the entire X-ray imaging spectrometer (XIS) bandpass (i.e. 0.6–10 keV in the observer frame) because the Fe K features are shifted to lower energies due to the high red-shift of the AGN.

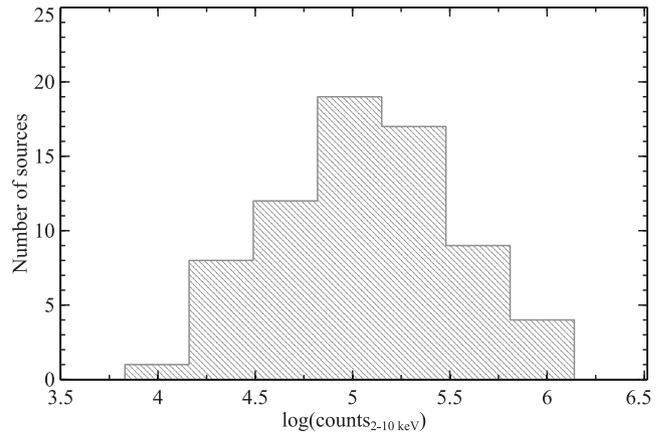
As highly ionized outflows are thought to originate at relatively small distances from the central nucleus (e.g. Gofford et al. 2011; Tombesi et al. 2011a, 2012, hereafter T11 and T12, respectively), it is important that the primary continuum emission from the central nucleus, rather than that which is reprocessed/scattered by circum-nuclear material out of line of sight, is directly observed so that such outflows can be detected. To this end we exclude all Type 2 sources to make sure that *all* sources were optically thin to X-rays below 10 keV. Therefore only those sources with a Type 1.0–1.9 orientation to the line of sight, as per the classifications listed on the *NASA/IPAC Extragalactic Database* (hereafter NED), in the VERONCAT catalogue itself and through literary sources were included. Note that we conservatively include the radio-quiet AGN ESO 103-G035 ( $z = 0.013\,29$ ) which, despite being classified as Type 2.0 in the VERONCAT and on NED, is often regarded as a Type 1.9 Seyfert in the literature (e.g. Warwick, Pounds & Turner 1988) by virtue of the presence of a moderately broad  $H\alpha$  line in its optical spectrum (Phillips et al. 1979). Observation details for all sources included in the heterogeneously selected sample are listed in Appendix A in the online supporting information. There are 99 observations of 51 AGN spanning a wide range of spectral types and radio-properties. As shown in Table 1 the sample is dominated by low–moderate inclination Seyfert 1.0–1.5 galaxies (28/51; 34/51 if Narrow Line Seyfert 1s are included) and contains comparatively few high-inclination (Type 1.8–1.9) Seyferts (6/51). There are six radio-loud sources in the sample, including all five of the BLRGs included in the T10B outflow case study, and five QSOs. The distributions of source red-shift and total 2–10 keV counts (for the XIS front-illuminated detector) are shown in Figs 1 and 2, respectively. The AGN are predominantly local, with  $\sim 90$  per cent of the sample being located at a red-shift of  $z \lesssim 0.1$ , but also includes more distant objects such as PDS 456 ( $z = 0.1840$ ), 1H 0419–577 ( $z =$

**Table 1.** Source classifications.

Source classification	Number
Sy 1.0–1.2	17
Sy 1.5	11
Sy 1.8–1.9	6
NLSy1	6
BLRG	6
QSO	5
Total	51



**Figure 1.** Histogram showing the distribution of sources included in this work. The high red-shift QSO APM 08279+5255 ( $z = 3.91$ ) has been omitted for scaling purposes.



**Figure 2.** Histogram showing the logarithm of the total front-illuminated XIS 2–10 keV counts in all fitted spectra. In stacked spectra (see Table A1) the total co-added counts are considered here, rather than the counts in each individual sequence.

0.1040), PKS 0558–504 ( $z = 0.1372$ ) and PBC J0839.7–1214 ( $z = 0.198$ ). The gravitationally lensed quasar APM 08279+5545 is by far the most distant object in the sample with a red-shift of  $z = 3.91$ . As shown in Fig. 2 the total 2–10 keV counts for the fitted spectra span two orders of magnitude but are approximately normally distributed about the peak and mean number of  $10^5$  counts.

### 3 DATA REDUCTION

#### 3.1 XIS reduction

There are four XIS CCD cameras aboard *Suzaku* which cover the energy range of 0.3–12.0 keV. The XIS 0, XIS 2 and XIS 3 are front-illuminated, while the XIS 1 is back-illuminated and provides superior sensitivity below  $\sim 1$  keV but a lower effective area and a higher background at harder X-ray energies. The XIS 2 suffered a charge leak in 2006 November<sup>3</sup> and is therefore only available in observations taken before this date. All spectra were reduced using

<sup>3</sup> <http://heasarc.gsfc.nasa.gov/docs/suzaku/news/xis2.html>

v6.11 of the *HEASoft*<sup>4</sup> suite following the process outlined in the *Suzaku* ABC Data Reduction Guide.<sup>5</sup> Events were selected using ASCA X-ray grades 0, 2, 3, 4 and 6, adopting the standard screening criteria such that data were excluded if taken: (1) within 436 s of passage through the South Atlantic Anomaly (SAA), (2) within an Earth elevation angle (ELV)  $<5^\circ$  and/or (3) with Earth day-time elevation angles  $<20^\circ$ . Hot and flickering pixels were removed from the XIS images using the `CLEANSIS` script. Source spectra were extracted from within circular regions of radius  $1.5 \leq r < 3.0$  arcmin to ensure good coverage of the source events, while background spectra were typically extracted from offset circles of the same radius with care taken to avoid the chip corners containing the Fe<sup>55</sup> calibration sources. For intrinsically faint sources the background spectra were extracted from circles larger than those of the source and the ratio of source/background area was accounted for with an appropriate background scaling factor. Redistribution matrices and ancillary response files for each observation were generated using the tasks `XISRMFGEN` and `XISSIMARFGEN`, respectively.

Where possible spectra obtained from the front-illuminated XIS detectors were combined into a single source spectrum (hereafter XIS-FI) using `MATHPHA` in order to maximize signal-to-noise ratio (S/N) in the Fe K band. Most of the observations in the sample have data for at least the XIS 0 and XIS 3 with a further 22 also having data for the XIS 2 (see Table A1). For SWIFT J2127.4+5654 (OB-SID 702122010) only data for the front-illuminated XIS 3 camera are available because the XIS 0 was not operational during the observation. Spectra for the back-illuminated XIS 1 (hereafter XIS-BI) were reduced using the same process as outlined above and were analysed simultaneously as a separate input spectrum. All XIS spectra were grouped to have a minimum of 50 counts per energy bin to enable the use of the  $\chi^2$  fit statistic. Net XIS exposures and the total background-subtracted 2–10 keV count rates (in the source rest-frame), for both the co-added XIS-FI and the XIS-BI detectors, are listed in Table A1.

### 3.2 HXD/PIN reduction

Spectra from the Hard X-ray Detector (HXD) Positive Intrinsic Negative (PIN) silicon diodes were also reduced using the method outlined in the *Suzaku* Data Reduction Guide and, again, processed according to the screening criteria described previously. As the HXD/PIN is a collimating rather than an imaging instrument, the contribution of both the instrumental non-X-ray background (NXB) and the cosmic X-ray background (CXB) needs to be independently accounted for when estimating the total background. The NXB was generated using the appropriate response and tuned event files (background model D) for each observation. Common good time intervals (GTIs) in the event and background files were selected using `MGTIME` and spectral files were extracted using `XSELECT`. PIN detector dead-time was accounted for using the `HXDDTCOR` task and the NXB background exposure was increased by ten times to reduce the effects of photon noise. The CXB contribution was simulated using the form of Boldt (1987), combined with the NXB to form a total background spectrum using `MATHPHA`, which was subsequently subtracted from the source spectrum within `XSPEC`.

All background-subtracted PIN spectra were binned to *at least* the  $3\sigma$  level above background (typically  $>5\sigma$ ) to enable the use of  $\chi^2$

fit statistic. Hard X-ray faint and/or high background observations where source count rates were  $<4$  per cent of the PIN total were not considered in our analysis. Only APM 08279+5255 and PDS 456, amounting to a total of five observations, meet this criterion (see Table A1). As with the XIS data, the final PIN exposures and the total background-subtracted rest-frame 15–50 keV counts are listed in Table A1.

## 4 ANALYSIS

### 4.1 Spectral fitting

A detailed broad-band spectral analysis of all sources was conducted using version 12.6.0q of the `XSPEC` spectral fitting package. All spectral models are modified by the appropriate column density of Galactic absorption using values taken from the Leiden/Argentine/Bonn Survey of Galactic H I (Kalberla et al. 2005) which were obtained from the on-line version of the `NH FTOOL`.<sup>6</sup> The values of Galactic absorption for each source are listed in Table A1. XIS-FI data were typically considered between 0.6 and 10.0 keV, while XIS-BI were only included between 0.6 and 5.0 keV due to decreasing S/N in the Fe K band which could hamper absorption line detection. All XIS spectra were excluded between 1.6 and 2.1 keV due to uncertainties in calibration associated with the Si K detector edge. Where available, PIN data were included to cover at least the 15–40 keV energy range.

A constant multiplicative factor was included in all models to account for the XIS/PIN cross-normalization whose value depends not only on the nominal pointing of the observation, but also the version of the *Suzaku* pipeline with which the data have been processed. The current XIS/PIN ratios suitable for version 2 processed data are 1.16 (1.18) for the XIS (HXD) nominal pointing positions.<sup>7</sup> However, the cross-normalization was up to  $\sim 5$ –6 per cent lower for data processed with version 1 of the pipeline, corresponding to an XIS/PIN ratio of 1.09 (1.13)<sup>7</sup> in the XIS (PIN) nominal pointing positions. While the difference is only small the additional uncertainty can have a considerable effect on the continuum parameters at  $E > 10$  keV, particularly in hard X-ray bright sources with a high PIN count rate where the model can become driven by the hard X-ray band. To account for any systematic effects associated with the instrumental cross-calibration we therefore allow the constant parameter in each model to vary  $\pm 5$  per cent about the values suggested by the *Suzaku* team to take into account any systematic errors.

There are 20 AGN in the sample which have two or more *Suzaku* observations. In APM 08279+5255, IC 4329A, MCG -6-30-15, Mrk 841, NGC 5506, NGC 5548 and PKS 0558–504 the different observations are similar in spectral shape which allows them to be co-added using the appropriate relative weighting factors to take into account differences in individual exposures, with the resultant time-averaged spectra for these sources being used in all subsequent analyses. 1H 0419–577 and NGC 2992 both had observations which were taken in different *Suzaku* nominal pointing positions which could influence co-adding. However, effective area was always consistent to within  $\pm 10$  per cent and so the spectra were still co-added using the mean of the response files. Any additional systematic uncertainty introduced was adequately accounted for by

<sup>4</sup> <http://heasarc.nasa.gov/lheasoft/>

<sup>5</sup> [ftp://legacy.gsfc.nasa.gov/suzaku/doc/suzaku\\_abc\\_guide.pdf](ftp://legacy.gsfc.nasa.gov/suzaku/doc/suzaku_abc_guide.pdf)

<sup>6</sup> <http://heasarc.gsfc.nasa.gov/cgi-bin/Tools/w3nh/w3nh.pl>

<sup>7</sup> <http://www.astro.isas.jaxa.jp/suzaku/doc/suzakumemo/suzakumemo-2007-11.pdf>

the variable XIS/PIN cross-normalization. In 3C 120 and Mrk 509, which both showed notable variability between observations, spectra were jointly fitted depending on the extent of their spectral variability between epochs. In 3C 120 we followed the analysis method of Kataoka et al. (2007) and T10B by co-adding OBSIDs 700001020, 700001030 and 700001040, which are all of a similar spectral shape and flux level, and jointly fitted them with OBSID 700001010 which has a more prominent underlying soft-excess. In Mrk 509, which is well known for having a strong soft-excess (Mehdipour et al. 2011), OBSID 701093010, the stacked OBSIDs 701093020, 701093030, 701093040 and OBSID 705025010 were fitted simultaneously to account for the observed variability in the soft X-ray band.

The remaining nine sources with more than one observation (3C 111, Fairall 9, Mrk 766, NGC 1365, NGC 3227, NGC 3516, NGC 3783, NGC 4051 and PDS 456) showed strong spectral variability in both the shape of the spectrum and/or drastic changes in flux state which made co-adding impractical. In these sources the available spectra were fitted simultaneously and a model was constructed to describe the observed spectral variability with as few additional free parameters as possible. In NGC 1365, OBSIDs 702047010 and 705031010 are dominated by very deep Fe K absorption lines. These lines are not present in OBSID 705031020, possibly due to the source dropping into a quasi-scattering-dominated state (see Maiolino et al. 2010 for details of the variability patterns in this source). For simplicity we therefore only simultaneously fitted OBSID 702047010 and OBSID 705031010 in NGC 1365, and fitted OBSID 705031020 separately.

The  $\chi^2$  minimization technique is used throughout this work, with all statistical errors quoted to the 90 per cent confidence level ( $\Delta\chi^2 = 2.71$  for one parameter of interest). Where the statistical significance of components is quoted in terms of a  $\Delta\chi^2$  value the component in question has been removed from the model and the data refitted to ensure that the order in which components are added to the model has no influence on the derived statistics. When referring to statistical changes to a fit a positive  $\Delta\chi^2$  denotes a worsening of the fit, while a negative  $\Delta\chi^2$  indicates a statistical improvement. Positive outflow velocities correspond to a net blue-shift relative to the systematic of the host galaxy, while a negative velocity indicates a net red-shift.

## 4.2 Model construction

All spectra were first fitted with a power law modified solely by Galactic absorption. Additional components were added to the model and retained provided their significance exceeded the >99 per cent confidence level by the  $F$ -test. Any emission lines in the soft X-ray band were fitted with narrow Gaussians ( $\sigma \sim 5$  eV). In sources with multiple spectra we initially found a broad-band fit to the observation with the highest flux; later observations were then added sequentially and the continuum/absorption parameters were allowed to vary independently until a simultaneous fit to all spectra had been achieved using as few free parameters as possible. Once a statistically acceptable fit to the broad-band spectrum of each source had been found, we fitted any necessary emission components and then performed a systematic and methodical search for Fe K absorption lines between 5 and 10 keV using energy–intensity plane contour plots (Section 4.3) and detailed Monte Carlo simulations (Section 4.4).

The modelling approach for individual spectral components are outlined below, with the continuum parameters for each source being noted in Table D1 and those for the warm absorber being

listed in Tables D2 and D3 for single- and multi-epoch spectra, respectively. A description of the underlying modelling assumptions and associated caveats is presented in Appendix C in the online supporting information.

### 4.2.1 Warm absorption I: full covering

Depending on the properties of the intervening material (such as its ionization state and column density) absorption by circumnuclear material can add significant spectral curvature to the observed X-ray spectrum and can therefore have a direct effect on the continuum and line parameters measured in broad-band models. In this work we model warm absorption components using a suite of *xstar* (v. 2.2.1bc) tables which are all generated assuming input values which are ‘typical’ for local Seyfert galaxies<sup>8</sup> (e.g. T11). The resultant *xstar* tables cover a wide range of parameter space in terms of column densities [ $10^{18} < \log(N_{\text{H}}/\text{cm}^{-2}) \leq 10^{24}$ ], and ionization parameter [ $-3 < \log(\xi/\text{erg cm s}^{-1}) \leq 6$ ] which makes them well suited for accounting for all manners of warm absorption.

Full-covering warm absorption zones are included in models where necessary to achieve a good fit to the soft X-ray band; in some cases more than one absorption zone is needed. In these cases the column density and ionization parameter of each zone are allowed to vary independently and represent an absorption geometry which consists of multiple layers of gas. At the energy resolution of the XIS CCDs the bound–bound absorption lines required to constrain the outflow velocities of individual soft X-ray absorption components are unresolved and all absorption zones are therefore fitted as stationary in the source rest-frame (i.e. fixed outflow velocities of  $v_{\text{out}} = 0$  km s<sup>-1</sup>). Allowing the outflow velocity of the soft X-ray absorber to vary always has a negligible effect on the reported Fe K absorption line parameters.

### 4.2.2 Warm absorption II: partial covering

We also consider the possibility that the sight line to a source is partially covered. In this absorption geometry a fraction  $f_{\text{cov}} < 1$  of the source flux is absorbed with the remaining  $1 - f_{\text{cov}}$  leaking through the absorption layer. For simplicity we account for partial-covering absorption layers using a customized version of the *zxipcf* model<sup>9</sup> which models the partial-covering absorption by partially ionized gas (see Reeves et al. 2008a) without needing to use complicated nested power laws and *xstar* tables. As in *zxipcf* the free parameters in our customized table model are column density ( $N_{\text{H}}$ ), ionization parameter ( $\log \xi$ ), covering fraction ( $f_{\text{cov}}$ ) and red-shift relative to the observer ( $z$ ). The model is based on the same tables as discussed in Section 4.2.1 and therefore has a slightly lower turbulent velocity than *zxipcf*, at  $v_{\text{turb}} = 100$  km s<sup>-1</sup>, but covers the same parameter space in terms of column density and ionization parameter.

Partial-covering absorption can have a strong effect on the observed continuum with moderate column densities of material ( $N_{\text{H}} \sim 10^{23}$  cm<sup>-2</sup>) adding considerable spectral curvature at  $E < 10$  keV (Reeves et al. 2004; Risaliti et al. 2005; Braito et al. 2011; Turner

<sup>8</sup> Absorption grids are described by an illuminating photon index of  $\Gamma = 2.1$ , a gas density of  $n = 10^{10}$  cm<sup>-3</sup>, a micro-turbulent velocity of  $v_{\text{turb}} = 100$  km s<sup>-1</sup> and an integrated model luminosity of  $L = 10^{44}$  erg s<sup>-1</sup> between 1 and 1000 Rydbergs.

<sup>9</sup> The standard *zxipcf* uses a specific grid with the following parameters:  $\Gamma = 2.2$ ,  $n = 10^{10}$  cm<sup>-3</sup>,  $v_{\text{turb}} = 200$  km s<sup>-1</sup>,  $L = 10^{44}$  erg s<sup>-1</sup>.

et al. 2011). Bound-free transitions in similar column density material can also fit broad residual emission profiles in the Fe K band (Inoue & Matsumoto 2003; Miller, Turner & Reeves 2008, 2009; Tatum et al. 2012), and partial covering by Compton-thick material ( $N_{\text{H}} \gtrsim 10^{24} \text{ cm}^{-2}$ ) can reduce the observed emission below 10 keV with the true intrinsic continuum only becoming apparent at higher energies as a ‘hard excess’ of emission relative to that expected from standard reflection models (Reeves et al. 2009; Risaliti et al. 2009; Turner et al. 2009; Tatum et al. 2013).

We include partially-covering absorbers in our models if and when they are required by the data at the  $P_{\text{F}} \geq 99$  per cent confidence level and a satisfactory fit to the data could not be achieved using solely fully-covering absorption; several sources appear to need multiple partially-covering absorption zones which suggests the presence of a clumpy stratified absorber along the line of sight. Again, all absorber parameters are listed in Tables D2 and D3 in Appendix D in the online supporting information. We note that in some circumstances the need of a high column density partially-covering component can be contingent on the means with which the underlying reflection component is modelled leading to some model degeneracies. However, as we discuss in Appendix C, regardless of whether the hard X-ray data are fitted with reflection or not, partial covering has little measurable effect on the parameters measured for any highly ionized absorption line systems.

#### 4.2.3 The soft-excess

Relative to the low energy extrapolation of the power-law continuum in the 2–10 keV band the X-ray spectrum of AGN often shows a smooth increase in emitted flux below  $\sim 1$  keV (Turner & Pounds 1988; Porquet et al. 2004). The thermal temperature of this ‘soft-excess’ suggests that it is unlikely to be the direct emission from a standard accretion disc without additional reprocessing (Sobolewska & Done 2007; Done et al. 2012). Alternative explanations posit that the soft-excess may be due to an increase in optical depth associated with circumnuclear O VII–VIII and Fe L-shell transitions at  $E \lesssim 0.7$  keV which can enhance either the transmitted or reflected flux along the sight line through smeared absorption (e.g. Chevallier et al. 2006; Done & Nayakshin 2007; Sobolewska & Done 2007) or blurred reflection (e.g. Crummy et al. 2006; Brenneman et al. 2011; Nardini et al. 2011; Nardini, Fabian & Walton 2012) effects. Furthermore, in some sources (e.g. Mrk 766, Miller et al. 2007; MCG -6-30-15, Miller et al. 2008) the ‘excess’ could simply be a product of complex absorption and just be the manifestation of opacity around  $\sim 1$ –2 keV.

Regardless of the true physical origin we take a purely phenomenological approach when fitting the soft-excess in this work, and predominantly use the `bbody` model which represents the emission from a constant temperature blackbody. While not necessarily physically motivated, modelling the soft-excess in this manner offers a simple parametrization of any ‘excess’ soft X-ray emission which is sufficient to get a good handle on the underlying continuum parameters. For completeness we investigate the effects that other ways of modelling the soft-excess can have on any Fe K-band absorption lines in Appendix C. Roughly half of the sources in the sample (24/51;  $\sim 47$  per cent) show evidence for a soft-excess, of which 22/24 are fitted with a `bbody` component. In 3C 120 and Mrk 509 the soft-excess is very broad and extends beyond that which can be fitted with a simple blackbody. In these sources we instead fit the excess with a second power law with a softer photon index.

#### 4.2.4 Lowly ionized reflection

Cold reflection from large column densities of neutral or lowly ionized material outside of the sight line can have a strong influence on the observed X-ray spectrum. The strongest observational characteristics of such reflection include the Compton-reflection hump at  $\sim 30$ –40 keV and the almost ubiquitous Fe  $K\alpha$  and Fe  $K\beta$  fluorescence lines at  $\sim 6.4$  and  $\sim 7.06$  keV (Nandra et al. 2007; Shu, Yaqoob & Wang 2010), respectively. Compton down-scattering of both  $K\alpha$  line photons and high energy continuum emission also gives rise to a ‘Compton shoulder’ at  $\sim 6.2$  keV (e.g. Matt 2002; Yaqoob & Murphy 2011) and resonant line emission in the soft X-ray band (e.g. Ross & Fabian 2005; García & Kallman 2010) which can further complicate the emergent spectrum.

Naturally, owing to the important effect it can have on the observed X-ray spectrum there are numerous models available for modelling the reflection component (e.g. `pexrav`/`pexriv`, Magdziarz & Zdziarski 1995; `pexmon`, Nandra et al. 2007; `reflionx`, Ross & Fabian 2005; `xillver`,<sup>10</sup> García & Kallman 2010; `MYTorus`,<sup>11</sup> Murphy & Yaqoob 2009). In this work we use a combination of `reflionx` and `pexrav`, both of which are publicly available and extensively used in the literature. The primary reason for this is that because `reflionx` interpolates the observed reflection spectrum from a pre-generated grid of table models it is significantly faster at fitting spectra than, for example, `pexrav`, `pexriv` or `pexmon`, which analytically calculate the Compton reflection spectrum on the fly at each step of the fitting process, or `MYTorus` which requires the model to be tailored for each individual source. Using `reflionx` as our primary means of fitting the reflection spectrum ensures the least time-consuming Monte Carlo simulations which is important when dealing with a large sample of objects such as that considered in this work. Secondly, when fitting reflection continuum without the simultaneous constraint of the Fe  $K\alpha$  line the reflection fraction,  $R$ , reported by `pexrav` (and `pexriv`) can become degenerate with the photon index of the primary power law, with a hardening reflection component compensated for by a softer primary continuum. By simultaneously fitting the reflection continuum, the Fe  $K\alpha$  line and the soft X-ray resonance lines, `reflionx` is able to overcome these modelling degeneracies which leads to a confident constraint on the contribution of the reflection continuum to the observed spectrum. `reflionx` also has the additional advantage of allowing the ionization state of the reflector,  $\xi$ , to be a free parameter which enables it to model changes in the Fe  $K\alpha$  emitted flux and Fe K-shell edge profile associated with the reflectors ionization state. We stress, however, that equivalent results are found for the detected Fe K absorbers if `pexmon` is used instead, but with the resultant Monte Carlo simulations taking significantly longer to complete which effectively prohibits its uniform use throughout the sample.

We initially fitted all sources with `pexrav` to determine the parameters of the Fe  $K\alpha$  and Fe  $K\beta$  lines. `pexrav` was then replaced with `reflionx`, and a systematic search for additional atomic features in the Fe K band was conducted. A total of 11 sources have best-fitting reflector Fe abundances which are non-solar, with 4 requiring a slight over-abundance (MCG+8-11-11, NCG 4593, NCG 7213, NCG 7469) and 7 with an under-abundance (4C +74.26, Fairall 9, IC 4329A, IGR J16185–5928, Mrk 335,

<sup>10</sup> The `xillver` reflection model is not currently available in the public domain and is therefore not considered for use in this work.

<sup>11</sup> The `MYTorus` model and documentation are publicly available at: [www.mytorus.com](http://www.mytorus.com)

Mrk 359, Mrk 841). These abundances are most likely a byproduct of `reflionx` assuming a face-on reflection geometry, and is a caveat which is discussed in greater detail in Appendix C. PDS 456 and APM 08279+5255 are not fitted with a reflection component because they lack sufficient counts in the `HXD/PIN`.

### 4.3 Searching for Fe K absorption

Once a statistically acceptable fit to the broad-band continuum had been found, i.e. including all necessary absorption regions, soft X-ray emission lines and continuum components, we performed a thorough search for additional spectral features between 5 and 10 keV. The method consists of inspecting the  $|\Delta\chi^2|$  deviations from the best-fitting continuum model using inverted contour plots of the energy–intensity plane in the Fe K band. The method of calculating the contour plots was adapted from the method outlined in T10A, and was carried out as follows:

(i) an unresolved ( $\sigma = 10$  eV) Gaussian was stepped across the entire 5–10 keV energy band of the baseline continuum model in 25 eV intervals, with normalization allowed to adopt both positive and negative values to probe for spectral lines in both emission and absorption. All of the other spectral components were allowed to be free;

(ii) after each step the  $\Delta\chi^2$  deviation was recorded generating a  $\chi^2$  distribution of the entire Fe K band relative to the baseline continuum model;

(iii) confidence contours for the grid of  $\chi^2$  values were plotted according to  $\Delta\chi^2$  deviations of  $-2.3$ ,  $-4.61$ ,  $-9.21$ ,  $-13.82$ ,  $-18.42$  from the baseline model, which correspond to confidence intervals of 68 per cent, 90 per cent, 99 per cent, 99.9 per cent and 99.99 per cent, respectively. A confidence contour corresponding to a  $\Delta\chi^2 = +0.5$  worse fit is also plotted which is intended to denote an approximate level for the continuum.

The energy–intensity contour plots produced with this method provide a powerful means of searching for additional emission or absorption components present in the Fe K band while also visually assessing their energy, intensity and rough statistical requirement relative to the underlying continuum model.

There are a number of atomic features between 6.0 and 8.0 keV which can complicate the detection of absorption line systems. Such features include ionized emission lines from Fe xxv and Fe xxvi expected at  $\sim 6.63$ – $6.7$  and  $\sim 6.97$  keV, respectively, the Fe K-shell edge complex at  $\sim 7.1$  keV, and any broad Fe line profile due to General Relativity (GR) or scattering effects. It is important to account for these emission residuals prior to searching for absorption lines, particularly in sources which show evidence for a broad residual as the broadness of the profile can effectively mask the presence of low-velocity absorption systems in the raw data. For a given continuum model we searched for highly ionized absorption lines in the Fe K band using the following steps:

(i) we first generated an energy–intensity contour plot using the method outlined above;

(ii) we then inspected the contour plot to determine whether there were any intense ionized emission and absorption lines present in the data with confidence contours of  $>99$  per cent;

(iii) where there was evidence for a broad emission residual at an  $F$ -test significance of at least  $P_F > 99$  per cent they were fitted with either a broadened Gaussian (with  $\sigma$ -width a free parameter) or a

`diskline`<sup>12</sup> depending on the asymmetry of the observed profile, and a second intermediate contour plot was generated to determine whether any further components were needed by the model. As before, all other model parameters were allowed to vary freely during this process. If there was no evidence for a broad profile we did not generate an additional contour plot and instead moved directly on to the next step;

(iv) where narrow emission profiles were detected with a resolved confidence contour of  $>99$  per cent they were fitted with unresolved ( $\sigma = 10$  eV fixed) Gaussian profiles provided they were required by the data at  $P_F > 99$  per cent (corresponding to  $\Delta\chi^2 > 9.21$  for two parameters of interest);

(v) once all emission profiles had been accounted for we again checked for the presence of any blue-shifted Fe K absorption lines at  $E > 6.7$  keV. If no absorption lines were present we ended the search at this step, reported the best-fitting continuum and Fe K emission line parameters in the relevant tables of Appendix D, and moved on to the next observation in the sample;

(vi) otherwise, where absorption troughs were clearly detected with a confidence contour of  $>99$  per cent we parametrized the line(s) using inverted Gaussians with  $\sigma$ -width initially fixed at either 10, 30 or 100 eV depending on which provided the greatest improvement to  $\Delta\chi^2$ . Note that while initially fixed, the line widths were allowed to vary where appropriate leading to four sources with resolved absorption profiles. The key parameters of all detected absorption lines are reported in Table 2.

This process was carried out on each individually fitted spectrum in the sample, including those which were included in simultaneous fits. In Figs 3 and 4 we show examples of this process applied to Ark 120 (OBSID 702014010), which shows evidence for a broad asymmetric emission profile and no absorption lines, and Mrk 766 (OBSID 701035020), which is dominated by Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  absorption. The top panels of both figures show the residuals which remain in the Fe K band when all atomic lines have been removed from the best-fitting continuum model and the reflection component is fitted with `pexrav` to highlight the presence of the neutral Fe K $\alpha$ /Fe K $\beta$  fluorescence lines. The contour plots show the significances of the remaining residuals when the Fe K $\alpha$  and Fe K $\beta$  lines have been fitted with `reflionx` and a narrow Gaussian line. The continuous outer contour corresponds to the  $\Delta\chi^2 = +0.5$  residual as mentioned previously. From outer to inner the closed significance contours corresponding to  $\Delta\chi^2$  improvements of  $-2.3$  ( $P_F = 68$  per cent),  $-4.61$  ( $P_F = 90$  per cent),  $-9.21$  ( $P_F = 99$  per cent),  $-13.82$  ( $P_F = 99.9$  per cent) and  $-18.42$  ( $P_F = 99.99$  per cent) relative to the best-fitting continuum model are shown in red, green, blue, cyan and magenta, respectively.

There are clear residual profiles present in both sources, particularly Ark 120 where the positive asymmetric profile and additional Fe xxvi emission are apparent in the middle and bottom panels of Fig. 3, respectively. The two absorption profiles in Mrk 766 – which have previously been identified as Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$ , respectively, by Miller et al. (2007) – are statistically distinguishable at  $>99.99$  per cent confidence and are self-consistently fitted with a single highly ionized region of photoionized absorption (see Section 4.6). Energy–intensity contour plots for all fitted spectra in

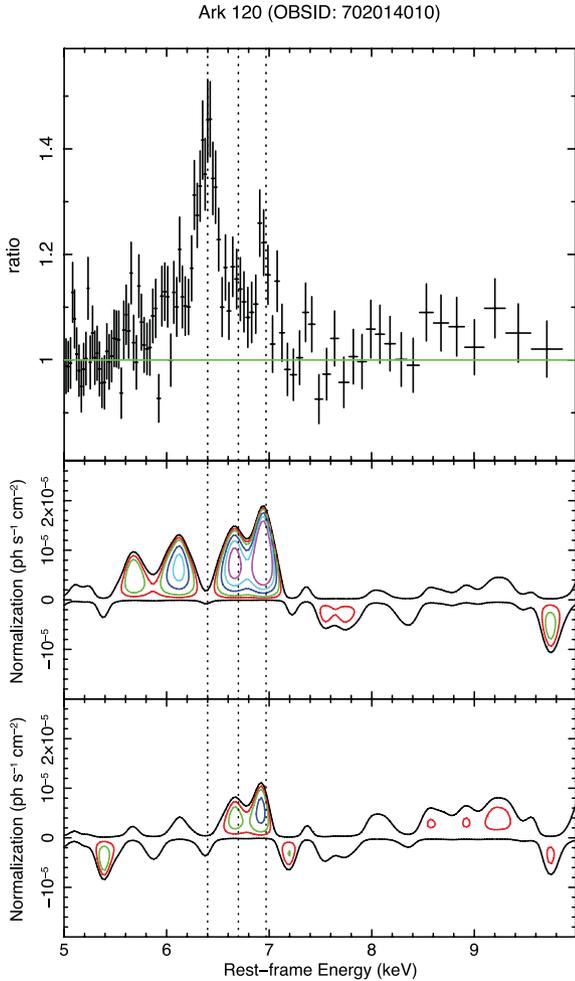
<sup>12</sup> For simplicity, the `diskline` was fitted with an assumed rest-frame energy of  $E = 6.4$  keV and a typical emissivity profile of  $q = -2.5$  (e.g. Patrick et al. 2011a). The inner emission radius ( $R_{in}$ ), outer radius ( $R_{out}$ ) and disc inclination ( $i$ ) were left free to vary.

**Table 2.** Gaussian parameters for detected Fe K band absorption lines. Notes: (1) source name; (2) *Suzaku* observation ID; (3) absorption line identification; (4) measured line energy in the source rest-frame, in units of keV; (5) measured line energy width, in units of eV. Unresolved lines were fitted with widths fixed at 10, 30 or 100 eV depending upon which yielded the best statistical improvement to the fit; (6) line intensity, in units of  $\times 10^{-6}$  photons  $\text{cm}^{-2} \text{s}^{-1}$ ; (7) line *EW*, in units of eV; (8) change in fit statistic (and degrees of freedom) when the line is removed from the best-fitting model; (9) line significance according to the *F*-test, in per cent; (10) Monte Carlo significance of individual lines, in per cent; (11) Monte Carlo significance of observing a pair of lines with energy separation corresponding to a Fe xxv–xxvi pair (see the text for details), in per cent. Absorption complexes detected at  $P_{\text{MC}} \geq 99$  per cent confidence are listed in bold.

Source (1)	OBSID (2)	Line ID (3)	Energy (4)	$\sigma$ -width (5)	Intensity (6)	<i>EW</i> (7)	$\Delta\chi^2/\Delta\nu$ (8)	$P_F$ (9)	$P_{\text{MC}}^1$ (10)	$P_{\text{MC}}^2$ (11)
3C 111	703034010	Ly $\alpha$	7.24 <sup>+0.04</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-5.6 <sup>+2.1</sup> <sub>-2.1</sub>	-26 <sup>+9</sup> <sub>-9</sub>	16.4/2	>99.99	<b>99.8</b>	
	705040020	Ly $\alpha$	7.76 <sup>+0.07</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-15.2 <sup>+5.1</sup> <sub>-5.1</sub>	-20 <sup>+10</sup> <sub>-10</sub>	11.1/2	99.60	97.8	
3C 390.3	701060010	Ly $\alpha$	8.08 <sup>+0.06</sup> <sub>-0.05</sub>	10 <sup>f</sup>	-5.6 <sup>+2.6</sup> <sub>-2.6</sub>	-21 <sup>+10</sup> <sub>-10</sub>	12.1/2	>99.99	98.7	
4+74.26	702057010	Ly $\alpha$	8.38 <sup>+0.05</sup> <sub>-0.07</sub>	10 <sup>f</sup>	-7.2 <sup>+3.4</sup> <sub>-3.4</sub>	-29 <sup>+14</sup> <sub>-14</sub>	12.4/2	99.80	96.6	
APM 08279+5255	stacked	He $\alpha$ -Ly $\alpha$ (1)	7.76 <sup>+0.11</sup> <sub>-0.11</sub>	100 <sup>f</sup>	-6.9 <sup>+2.5</sup> <sub>-2.3</sub>	-24 <sup>+8</sup> <sub>-8</sub>	23.8/2	>99.99	<b>99.7</b>	
		He $\alpha$ -Ly $\alpha$ (2)	10.97 <sup>+0.16</sup> <sub>-0.16</sub>	100 <sup>f</sup>	-4.4 <sup>+1.9</sup> <sub>-1.8</sub>	-30 <sup>+15</sup> <sub>-15</sub>	18.4/2	>99.99	<b>97.3</b>	
CBS 126	705042010	Ly $\alpha$	7.04 <sup>+0.04</sup> <sub>-0.05</sub>	30 <sup>f</sup>	-3.3 <sup>+1.0</sup> <sub>-1.0</sub>	-77 <sup>+22</sup> <sub>-22</sub>	27.5/2	>99.99	> <b>99.9</b>	
ESO 103-G035	703031010	Ly $\alpha$	7.26 <sup>+0.05</sup> <sub>-0.05</sub>	104 <sup>+61</sup> <sub>-55</sub>	-14.7 <sup>+4.8</sup> <sub>-4.1</sub>	-40 <sup>+14</sup> <sub>-14</sub>	35.0/3	>99.99	> <b>99.9</b>	
MCG -6-30-15	stacked[all]	He $\alpha$	6.72 <sup>+0.02</sup> <sub>-0.01</sub>	10 <sup>f</sup>	-7.8 <sup>+1.2</sup> <sub>-1.2</sub>	-18 <sup>+3</sup> <sub>-3</sub>	81.6/2	>99.99	> <b>99.9</b>	> <b>99.9</b>
		Ly $\alpha$	7.05 <sup>+0.02</sup> <sub>-0.02</sub>	10 <sup>f</sup>	-5.1 <sup>+1.2</sup> <sub>-1.2</sub>	-14 <sup>+3</sup> <sub>-3</sub>	25.7/2	>99.99	> <b>99.9</b>	
MR 2251-178	704055010	Fe L-He $\alpha$	7.55 <sup>+0.12</sup> <sub>-0.12</sub>	193 <sup>+84</sup> <sub>-62</sub>	-12.9 <sup>+4.1</sup> <sub>-4.1</sub>	-32 <sup>+10</sup> <sub>-10</sub>	26.9/3	>99.99	<b>99.6</b>	
Mrk 279	704031010	Ly $\alpha$	8.88 <sup>+0.06</sup> <sub>-0.06</sub>	10 <sup>f</sup>	-2.5 <sup>+1.0</sup> <sub>-1.0</sub>	-73 <sup>+30</sup> <sub>-30</sub>	15.1/2	99.95	<b>99.1</b>	
Mrk 766	701035010	He $\alpha$	7.11 <sup>+0.07</sup> <sub>-0.07</sub>	10 <sup>f</sup>	-2.8 <sup>+1.6</sup> <sub>-1.6</sub>	-27 <sup>+15</sup> <sub>-15</sub>	7.0/2	96.27	71.0	<b>99.8</b>
		Ly $\alpha$	7.42 <sup>+0.06</sup> <sub>-0.06</sub>	10 <sup>f</sup>	-3.9 <sup>+1.6</sup> <sub>-1.6</sub>	-40 <sup>+17</sup> <sub>-17</sub>	15.3/2	99.92	98.4	
	701035020	He $\alpha$	6.80 <sup>+0.02</sup> <sub>-0.02</sub>	10 <sup>f</sup>	-7.7 <sup>+1.8</sup> <sub>-1.8</sub>	-60 <sup>+14</sup> <sub>-14</sub>	50.0/2	>99.99	> <b>99.9</b>	> <b>99.9</b>
		Ly $\alpha$	7.12 <sup>+0.04</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-5.0 <sup>+1.9</sup> <sub>-1.9</sub>	-44 <sup>+17</sup> <sub>-17</sub>	18.9/2	>99.99	<b>99.8</b>	
NGC 1365	702047010	He $\alpha$	6.77 <sup>+0.01</sup> <sub>-0.01</sub>	65 <sup>+8</sup> <sub>-8</sub>	-29.9 <sup>+1.6</sup> <sub>-1.6</sub>	-156 <sup>+8</sup> <sub>-8</sub>	1123.8/3	>99.99	> <b>99.9</b>	> <b>99.9</b>
		Ly $\alpha$	6.97 <sup>+0.03</sup> <sub>-0.03</sub>	65*	-8.7 <sup>+1.7</sup> <sub>-1.7</sub>	-69 <sup>+13</sup> <sub>-13</sub>	77.0/3	>99.99	> <b>99.9</b>	
		He $\beta$	7.94 <sup>+0.01</sup> <sub>-0.01</sub>	65*	-26.1 <sup>+1.5</sup> <sub>-1.5</sub>	-164 <sup>+10</sup> <sub>-10</sub>	964.5/3	>99.99	> <b>99.9</b>	
		Ly $\beta$	8.38 <sup>+0.03</sup> <sub>-0.03</sub>	65*	-10.6 <sup>+1.8</sup> <sub>-1.8</sub>	-87 <sup>+15</sup> <sub>-15</sub>	94.3/3	>99.99	> <b>99.9</b>	
	704031010	He $\alpha$	6.71 <sup>+0.02</sup> <sub>-0.02</sub>	30 <sup>f</sup>	-16.1 <sup>+2.4</sup> <sub>-2.4</sub>	-67 <sup>+10</sup> <sub>-10</sub>	134.9/2	>99.99	> <b>99.9</b>	> <b>99.9</b>
		Ly $\alpha$	7.00 <sup>+0.02</sup> <sub>-0.02</sub>	30*	-13.4 <sup>+2.0</sup> <sub>-2.0</sub>	-85 <sup>+14</sup> <sub>-14</sub>	90.7/2	>99.99	> <b>99.9</b>	
		He $\beta$	7.87 <sup>+0.08</sup> <sub>-0.08</sub>	30*	-4.7 <sup>+2.5</sup> <sub>-2.5</sub>	-34 <sup>+20</sup> <sub>-20</sub>	7.2/2	96.77	-	
		Ly $\beta$	8.37 <sup>+0.11</sup> <sub>-0.11</sub>	30*	-5.3 <sup>+2.7</sup> <sub>-2.7</sub>	-45 <sup>+24</sup> <sub>-24</sub>	8.5/2	98.23	-	
NGC 3227	703022010	He $\alpha$	6.69 <sup>+0.04</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-9.9 <sup>+3.7</sup> <sub>-3.7</sub>	-21 <sup>+7</sup> <sub>-7</sub>	17.5/2	99.99	> <b>99.9</b>	> <b>99.9</b>
		Ly $\alpha$	6.95 <sup>+0.03</sup> <sub>-0.03</sub>	10 <sup>f</sup>	-15.5 <sup>+3.7</sup> <sub>-3.7</sub>	-38 <sup>+9</sup> <sub>-9</sub>	42.4/2	>99.99	> <b>99.9</b>	
	703022030	He $\alpha$	6.76 <sup>+0.03</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-8.5 <sup>+3.5</sup> <sub>-3.5</sub>	-24 <sup>+9</sup> <sub>-9</sub>	16.2/2	99.94	<b>99.8</b>	> <b>99.9</b>
		Ly $\alpha$	7.04 <sup>+0.05</sup> <sub>-0.05</sub>	10 <sup>f</sup>	-6.6 <sup>+3.6</sup> <sub>-3.6</sub>	-22 <sup>+11</sup> <sub>-11</sub>	9.3/3	98.64	91.8	
703022050	He $\alpha$	6.76 <sup>+0.07</sup> <sub>-0.07</sub>	10 <sup>f</sup>	-5.8 <sup>+2.8</sup> <sub>-2.8</sub>	-19 <sup>+9</sup> <sub>-8</sub>	12.0/2	99.74	96.2	> <b>99.9</b>	
	Ly $\alpha$	7.05 <sup>+0.03</sup> <sub>-0.03</sub>	10 <sup>f</sup>	-11.0 <sup>+2.8</sup> <sub>-2.8</sub>	-41 <sup>+9</sup> <sub>-9</sub>	41.2/2	>99.99	> <b>99.9</b>		
NGC 3516	100031010	He $\alpha$	6.74 <sup>+0.02</sup> <sub>-0.02</sub>	10 <sup>f</sup>	-9.1 <sup>+1.6</sup> <sub>-1.6</sub>	-25 <sup>+5</sup> <sub>-5</sub>	65.6/2	>99.99	> <b>99.9</b>	> <b>99.9</b>
		Ly $\alpha$	7.00 <sup>+0.02</sup> <sub>-0.02</sub>	10 <sup>f</sup>	-6.0 <sup>+1.7</sup> <sub>-1.7</sub>	-20 <sup>+6</sup> <sub>-6</sub>	66.7/2	>99.99	> <b>99.9</b>	
NGC 3783	701033010	He $\alpha$	6.69 <sup>+0.04</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-10.0 <sup>+3.2</sup> <sub>-3.0</sub>	-16 <sup>+5</sup> <sub>-6</sub>	30.2/2	>99.99	> <b>99.9</b>	
	704063010	He $\alpha$	6.71 <sup>+0.02</sup> <sub>-0.02</sub>	10 <sup>f</sup>	-14.0 <sup>+3.1</sup> <sub>-3.1</sub>	-20 <sup>+3</sup> <sub>-5</sub>	56.7/2	>99.99	> <b>99.9</b>	
NGC 4051	700004010	He $\alpha$	6.77 <sup>+0.04</sup> <sub>-0.04</sub>	10 <sup>f</sup>	-2.2 <sup>+1.0</sup> <sub>-1.0</sub>	-21 <sup>+10</sup> <sub>-10</sub>	9.9/2	99.39	89.4	> <b>99.9</b>
		Ly $\alpha$	7.08 <sup>+0.03</sup> <sub>-0.03</sub>	10 <sup>f</sup>	-3.3 <sup>+1.1</sup> <sub>-1.1</sub>	-38 <sup>+13</sup> <sub>-13</sub>	22.7/2	>99.99	> <b>99.9</b>	
703023010	He $\alpha$	6.82 <sup>+0.03</sup> <sub>-0.03</sub>	10 <sup>f</sup>	-4.8 <sup>+1.3</sup> <sub>-1.3</sub>	-22 <sup>+6</sup> <sub>-6</sub>	26.4/2	>99.99	> <b>99.9</b>	> <b>99.9</b>	
	Ly $\alpha$	7.11 <sup>+0.02</sup> <sub>-0.02</sub>	10 <sup>f</sup>	-6.4 <sup>+1.3</sup> <sub>-1.3</sub>	-33 <sup>+7</sup> <sub>-7</sub>	57.5/2	>99.99	> <b>99.9</b>		
NGC 4151	701034010	Ly $\alpha$	7.17 <sup>+0.04</sup> <sub>-0.04</sub>	92 <sup>+51</sup> <sub>-41</sub>	-15.9 <sup>+4.7</sup> <sub>-4.3</sub>	-28 <sup>+8</sup> <sub>-8</sub>	43.7/3	>99.99	> <b>99.9</b>	
NGC 4395	702001010	He $\alpha$	6.63 <sup>+0.07</sup> <sub>-0.05</sub>	10 <sup>f</sup>	-2.2 <sup>+1.1</sup> <sub>-1.1</sub>	-35 <sup>+19</sup> <sub>-19</sub>	10.4/2	>99.99	74.5	> <b>99.9</b>
		Ly $\alpha$	6.91 <sup>+0.06</sup> <sub>-0.06</sub>	10 <sup>f</sup>	-2.9 <sup>+1.1</sup> <sub>-1.1</sub>	-55 <sup>+20</sup> <sub>-20</sub>	18.3/2	>99.99	86.8	
NGC 5506	stacked[all]	Ly $\alpha$	9.23 <sup>+0.06</sup> <sub>-0.06</sub>	10 <sup>f</sup>	-11.1 <sup>+4.7</sup> <sub>-4.7</sub>	-16 <sup>+5</sup> <sub>-5</sub>	16.2/2	>99.99	<b>99.8</b>	
PDS 456	701056010	He $\alpha$ -Ly $\alpha$ (1)	9.07 <sup>+0.06</sup> <sub>-0.06</sub>	100 <sup>f</sup>	-2.5 <sup>+0.8</sup> <sub>-0.8</sub>	-108 <sup>+35</sup> <sub>-35</sub>	29.5/2	>99.99	> <b>99.9</b>	
		He $\alpha$ -Ly $\alpha$ (2)	9.57 <sup>+0.09</sup> <sub>-0.09</sub>	100 <sup>f</sup>	-2.0 <sup>+0.8</sup> <sub>-0.8</sub>	-99 <sup>+41</sup> <sub>-41</sub>	15.7/2	>99.99	<b>99.8</b>	
	705041010	He $\alpha$ -Ly $\alpha$ (1)	8.58 <sup>+0.09</sup> <sub>-0.09</sub>	100 <sup>f</sup>	-2.8 <sup>+1.0</sup> <sub>-1.0</sub>	-118 <sup>+46</sup> <sub>-46</sub>	19.5/2	>99.99	96.2	
		He $\alpha$ -Ly $\alpha$ (2)	9.03 <sup>+0.09</sup> <sub>-0.09</sub>	100 <sup>f</sup>	-3.5 <sup>+1.1</sup> <sub>-1.1</sub>	-169 <sup>+51</sup> <sub>-51</sub>	78.6/2	>99.99	> <b>99.9</b>	
S J2127.4+5654	702122010	Ly $\alpha$	9.04 <sup>+0.05</sup> <sub>-0.05</sub>	10 <sup>f</sup>	-10.0 <sup>+4.7</sup> <sub>-4.7</sub>	-60 <sup>+28</sup> <sub>-28</sub>	11.9/2	99.78	98.9	

<sup>f</sup> indicates that a parameter was frozen during spectral fitting.

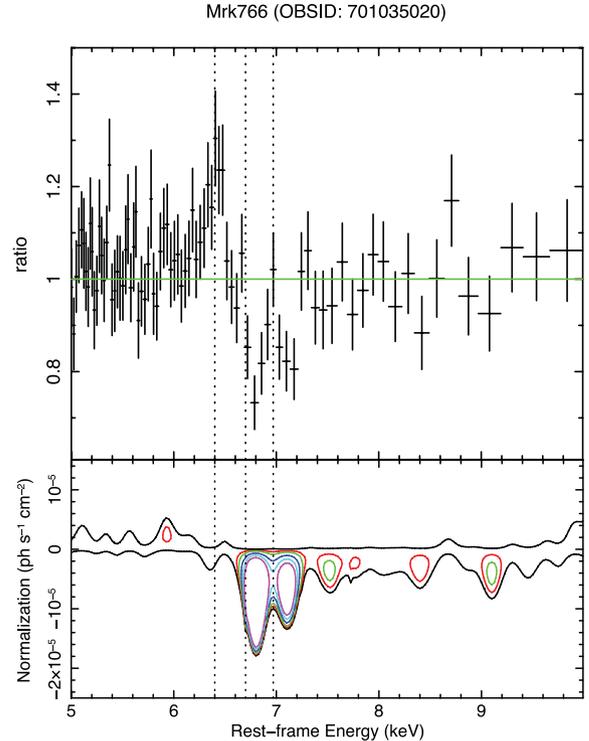
\* indicates that a parameter was tied during spectral fitting.



**Figure 3.** Top panel: ratio plot showing the residuals remaining in the Fe K band of Ark 120 (OBSID 702014010) once all atomic lines have been removed and reflection fitted with `pxravn`. Middle panel: confidence contours showing deviations from the best-fitting model when the Fe K $\alpha$  and Fe K $\beta$  lines have been fitted with `reflionx` and a narrow Gaussian, respectively. The closed significance contours corresponding (from outer to inner) to  $\Delta\chi^2$  improvements of  $-2.3$  ( $P_F = 68$  per cent),  $-4.61$  ( $P_F = 90$  per cent),  $-9.21$  ( $P_F = 99$  per cent),  $-13.82$  ( $P_F = 99.9$  per cent) and  $-18.42$  ( $P_F = 99.99$  per cent) relative to the best-fitting continuum model are shown in red, green, blue, cyan and magenta, respectively. The magenta contours in the middle panel indicate that the broad emission residual is significant at the  $>99.99$  per cent level and that an additional component is required in the model. Bottom panel: the remaining confidence contours once the broad profile has been fitted. The broad residual are no longer detected but an additional Fe xxvi emission line is present at  $E \sim 6.97$  keV which is significant at  $>99$  per cent. In all panels the dashed vertical lines indicate the expected rest-frame energies of (from left to right) the Fe K $\alpha$  fluorescence line, Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$ , respectively.

this sample are included in Figs B1 and B2 in Appendix B, in the online supporting information. We note that the  $\Delta\chi^2$  of individual lines and their corresponding significances according to the  $F$ -test as listed in Table 2 are taken from a direct spectral fitting of the line profiles themselves, and are not determined directly from the contour plots.

It is also important to note that while we use the  $F$ -test as a rough initial gauger of the significance of any line-like profiles in the Fe K band we do not claim to robustly detect any absorption lines based solely on their measured  $\Delta\chi^2$  and corresponding  $F$ -test



**Figure 4.** Top panel: as in the case of Fig. 3, but this time for Mrk 766 (OBSID 701035020). Bottom panel: as in the corresponding panel of Fig. 3. There are no ionized emission lines required in Mrk 766; however, two highly significant (each required at  $P_F > 99.99$  per cent) absorption profiles are clearly detected. The energy of these lines is consistent with Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$ , respectively, outflowing at  $v_{\text{out}} \sim 6000$  km s $^{-1}$ .

significances. Indeed, as has been pointed out by numerous authors in the literature (e.g. Protassov et al. 2002; Porquet et al. 2004; Markowitz, Reeves & Braito 2006; T10A) the  $F$ -test alone might not be an adequate statistical test to determine the detection significance of atomic lines in complex spectral models as it neither takes into account the number of energy resolution bins over a given energy range, nor the expected energy of a given atomic line. In cases where there is no a priori justification for expecting a spectral line at a particular energy and the line search is done over what is essentially an arbitrary energy range, as is often the case when dealing with strongly blue-shifted absorption lines, the  $F$ -test can somewhat overpredict the detection probability and when compared to extensive simulations. For these reasons all suspected absorption lines which have an  $F$ -test significance of  $P_F > 99$  per cent were followed up with extensive Monte Carlo simulations which allows their detection significance to be assessed against random fluctuations and photon noise in the spectrum.

#### 4.4 Monte Carlo simulations

Such Monte Carlo simulations have been used extensively in the literature to overcome the limitations of the  $F$ -test (e.g. Porquet et al. 2004; Markowitz et al. 2006; Miniutti & Fabian 2006; T10A; T10B) and enable the statistical significance of a spectral line to be robustly determined independently of spectral noise and detector effects. The method of Monte Carlo simulation we used follows the same process which was first outlined by Porquet et al. (2004), and is almost identical to that used by T10A. The process was carried out with the following steps.

(i) From the null hypothesis model (i.e. the broad-band continuum model with all narrow absorption lines in the Fe K band removed) we simulated both XIS-FI and XIS-BI spectra using the *fakeit* command in *XSPEC* and subtracted the appropriate background files. The simulated spectra had the same exposure as the original data and used the same spectral response files.

(ii) We then fitted the simulated XIS-FI(XIS-BI) spectra between 0.6 and 10.0 keV (0.6–5.0 keV) with the null hypothesis model to produce a new and refined null hypothesis which takes into account the uncertainty in the null hypothesis model itself. All continuum parameters bar the photon index of the primary power law and its normalization, and the normalization of any *tbody* component, were frozen to their best-fitting parameters taken from the real data to prevent degeneracies between model components during re-fitting. Any broad profile at Fe K had its width fixed to the best-fitting value but was allowed to vary in both centroid energy and normalization.

(iii) From the refined null hypothesis model we generated a second set of simulated XIS-FI and XIS-BI spectra and subtracted the appropriate background files. These second simulated spectra were then fitted with the null hypothesis model and the resultant  $\chi^2_{\text{null}}$  value was recorded.

(iv) An unresolved ( $\sigma = 10$  eV) Gaussian was added to the model at 5 keV in the source rest-frame with intensity initially set to zero but left free to vary between both positive and negative values to probe for both emission and absorption lines. The Gaussian line was then sequentially stepped between 5 and 9.5 keV (rest-frame) in 25 eV increments. After each step the  $|\Delta\chi^2_{\text{noise}}|$  was recorded relative to  $\chi^2_{\text{null}}$ .

(v) This process was carried out for  $T = 1000$  simulated spectra per observation yielding a distribution of  $|\Delta\chi^2_{\text{noise}}|$  under the null hypothesis which essentially maps the statistical significance of any deviations from the null hypothesis model which are solely due to random photon noise in the spectrum.

(vi) The measured significance of the line in the real spectrum  $|\Delta\chi^2_{\text{line}}|$  was then compared to the values in the  $|\Delta\chi^2_{\text{noise}}|$  distribution to assess how many simulated spectra had a random fluctuation with a detection significance over this threshold value. If  $N$  simulated spectra have  $|\Delta\chi^2_{\text{noise}}| \geq |\Delta\chi^2_{\text{line}}|$ , then the estimated confidence interval for the observed line from Monte Carlo simulations is then  $P_{\text{MC}}^{\text{Line}} = 1 - \left(\frac{N_{\text{Line}}}{T}\right)$ . Moreover, if there are two absorption profiles consistent with a Fe  $\text{XXV}$ – $\text{XXVI}$  pair, we can infer the null probability of both lines simultaneously being a false detection by multiplying the probabilities of each individual line.

Monte Carlo significances for all absorption lines detected in this work are listed in column 10 of Table 2. Absorption complexes with a total  $P_{\text{MC}} \geq 99$  per cent are conservatively identified as robustly detected, while those with 95 percent  $\leq P_{\text{MC}} < 99$  per cent are only listed as marginal detections. The overall detection rate and global significance of Fe K-band absorption lines are further discussed in Section 5.1.

There are two possible caveats associated with the above Monte Carlo process which both warrant further discussion. First and foremost, the Monte Carlo simulations implicitly assume that the null-hypothesis model is the correct representation of the continuum in a given source, and therefore the Monte Carlo probabilities do not account for the possibility of continuum mis-modelling. Given this possibility it is important to note that we have attempted to achieve a statistically acceptable representation of the broad-band spectrum, so that no obvious broad residuals are present. Care has been taken in the Fe K band in particular such that any broad residuals are

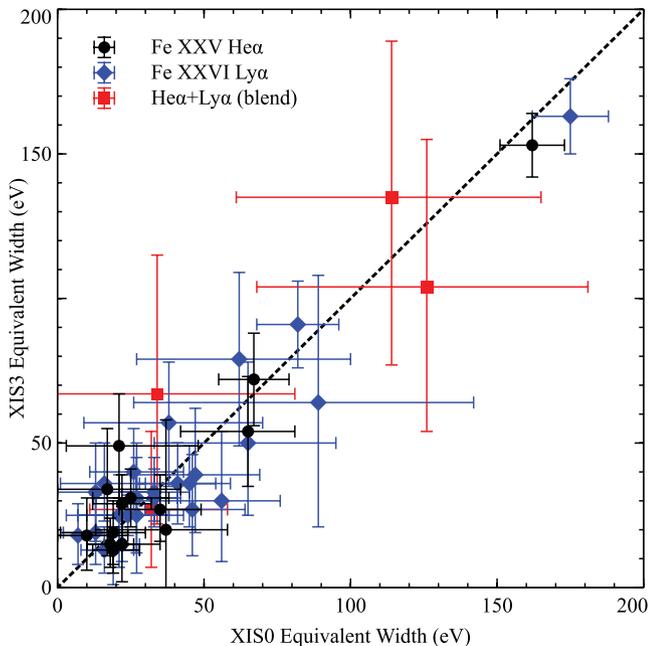
minimized prior to searching for absorption lines, such that the reduced  $\chi^2$  is  $\approx 1.0$  in all cases. Moreover, and as shown in Appendix, we note that the vast majority of suspected absorption residuals are manifested through discrete narrow dips in the spectrum relative to the best-fitting model, whereas any systematic residuals would usually be broader than the instrument resolution.

The second caveat is associated with spectral complexity in the Fe K band which can complicate both line identification and spectral interpretation. There can be significant atomic complexity between  $\sim 5$  and 7 keV, e.g. the narrow Fe  $K\alpha$  and Fe  $K\beta$  fluorescence lines, broad underlying Fe K lines and ionized Fe  $\text{XXV}$ – $\text{XXVI}$  emission lines, and the detection of absorption in this regime can depend strongly on how these features are modelled. There are three sources (4C +74.26, MGC-6-30-15, SWIFT J2127.4+5654; see Fig. B2) in the sample in which both a broad underlying profile and at least one absorption line have been detected. In both 4C+74.26 and SWIFT J2127.4+5654 the absorption lines are detected at high energies (i.e.  $E > 8$  keV in the source rest-frame) and the effect of the broad Fe line on the line detection is negligible. In MGC -6-30-15 the lines have also been confirmed with other X-ray observatories (Young et al. 2005; Miller et al. 2009) which suggests that the presence of a broad Fe line does not introduce any significant model systematics. Thus modelling of a broad emission line does not appear to affect the detections of Fe K absorption lines in these cases.

#### 4.5 Consistency checks

To further test the robustness of the absorption lines detected in the co-added XIS-FI spectra we performed a series of consistency checks with the individual XIS detectors. If a line is detected in two (or more, if the XIS 2 is also present) detectors, the line is very likely to be a real feature intrinsic to the source spectra rather than an artefact of background subtraction, systematic noise or an associated detector effect. The XIS-BI is not suitable for use as a consistency check as it tends to have a much lower S/N above around 5 keV owing to its lower effective area and higher background. Where absorption has been detected at the  $P_{\text{MC}} \geq 95$  per cent level we fitted each individual background-subtracted XIS 0, XIS 2 (where present) and XIS 3 spectrum with the best-fitting model to the co-added XIS-FI spectrum. We then fitted a Gaussian absorption line, with rest-frame energy and normalization left free to vary, at the energy where the absorption line is detected in the XIS-FI spectrum, refitted the data and noted the resultant line parameters for each XIS detector. We note that in SWIFT J2127.4+5654 this consistency check could not be conducted as only the XIS 3 spectrum was available.

In all sources the residuals detected in the individual XIS spectra have Gaussian parameters consistent with those found for the co-added FI spectrum. In Fig. 5 we show a comparison of absorption line equivalent width ( $EW$ ), along with the 90 per cent error bars, as measured independently with the XIS 0 and XIS 3 detectors. The dashed diagonal line represents the position where the  $EW$  is equal in each of the detectors; where the  $EW$  is consistent at the 90 per cent level the error bars will overlap this line. In this plot the Fe  $\text{XXV}$   $\text{He}\alpha$ , Fe  $\text{XXVI}$   $\text{Ly}\alpha$  and blended Fe  $\text{XXV}$ – $\text{XXVI}$  lines are plotted separately, even in the cases where the lines are found to be part of a Fe  $\text{XXV}$  pair. We find that the mean  $EW$ s measured with the XIS 0 and XIS 3 detectors are consistent at the 90 per cent level (i.e. the error bars cross the diagonal line in Fig. 5) in all sources. Where possible this analysis was also carried out using the XIS 2 spectra versus both XIS 0 and XIS 3. Again, in all cases the mean parameters of the lines were always consistent at the 90 per cent



**Figure 5.** Consistency check comparing the absorption line *EW* measured with the XIS 0 versus the XIS 3 along with the 90 percent error bars. The diagonal dashed line represents the location where the *EW* is equal in both detectors. The circle (black), diamond (blue) and square (red) markers show the *EW* of the Fe xxv He $\alpha$ , Fe xxvi Ly $\alpha$  and blended Fe xxv–xxvi lines, respectively. For sources with lines detected in more than one observation the mean *EW* for each line in each detector and associated errors are plotted. Where available this consistency check was also carried out with the XIS 2 which yielded results which are entirely consistent with those displayed here.

level. The consistency check strongly suggests that the observed lines are real features and are not due to detector effects or due to background subtraction.

#### 4.6 Photoionization modelling

All absorption complexes detected at a  $P_{MC} \geq 95$  percent confidence level were then fitted with the *xstar* photoionization code to probe the likely physical properties of the absorbing material, such as the column density  $N_H$ , ionization parameter  $\xi$  and the red-shift of the absorber relative to the observer  $z_o$ . Probing these parameters is important as not only do they allow the mean properties of the absorbing material to be determined but, through the use of simple geometric assumptions, they also permit an order-of-magnitude assessment of an absorber likely radial distance from the ionizing source, the mass outflow rate and their global energetic output. Detailed discussion regarding the absorber kinematics will be presented in a companion paper (Gofford et al., in preparation).

In this work the absorption lines were fitted with an *xstar* table that had an assumed illuminating continuum of  $\Gamma = 2.1$  and a micro-turbulent velocity ( $v_{\text{turb}}$ ) of  $1000 \text{ km s}^{-1}$ , which roughly corresponds to the full width at half-maximum (FWHM) velocity width of a  $\sigma = 10 \text{ eV}$  Fe xxvi Ly $\alpha$  absorption line at a rest-frame energy of  $6.97 \text{ keV}$ . In cases where the lines were resolved (i.e.  $\sigma > 10 \text{ eV}$ ) we used *xstar* tables with  $v_{\text{turb}}$  values which closely matched the measured FWHM velocity width of the observed profile we used; we found that tables with  $v_{\text{turb}}$  equal to  $1000, 3000, 5000$  and  $10000 \text{ km s}^{-1}$  were sufficient to fit all systems in the sample. In 17/20 sources with significantly detected absorption lines all of the parameters of

the *xstar* table were allowed to vary freely. In APM 08279+5255 and the joint spectra fits to NGC 1365 and PDS 456 some parameters were tied to prevent degeneracies between the column density and ionization parameter (see Table 3). The absorber red-shift ( $z_o$ ) as measured from the spectrum using *xstar* is given in the observer frame. This is related to the intrinsic absorber red-shift in the source rest-frame ( $z_a$ ) and the cosmological red-shift of the source ( $z_c$ ) through the relation  $(1 + z_o) = (1 + z_a)(1 + z_c)$ . From this, the intrinsic velocity of the outflow relative to the source ( $v_{\text{out}}$ ) can then be calculated from the relativistic Doppler formula which ensures that the relativistic effects associated with both high red-shift sources and high-velocity outflows are correctly taken into account when inferring absorber outflow velocities relative to the source rest-frame.

##### 4.6.1 Line identifications

Before discussing the results of the *xstar* fitting it is important that the a priori assumption that the absorption lines detected at  $E \gtrsim 6.6 \text{ keV}$  are due to the velocity shifted resonance lines of Fe xxv and Fe xxvi is justified. Indeed, while the K $\alpha$  transitions of Fe xxv and Fe xxvi, which are expected at mean rest-frame energies of  $6.697$  and  $6.966 \text{ keV}$ , respectively, are expected to be the strongest lines in the  $\approx 6.5\text{--}7.0 \text{ keV}$  energy interval, there are several other atomic features at higher energies which may complicate the identification of blue-shifted absorption systems at  $E > 7 \text{ keV}$ . For example, the K-shell edges from the various ionized species of Fe are found above  $7 \text{ keV}$ , with energies ranging from  $7.1 \text{ keV}$  for neutral Fe up to  $9.3 \text{ keV}$  for Fe xxvi. In the case of low-moderate ionization Fe the K-edge is accompanied by higher order resonance line structure which can give the edge a subtle curved profile in CCD spectra rather than it simply being an abrupt drop in flux (Kallman et al. 2004). Furthermore, given its proximity to the rest-frame energy of the Fe xxv–xxvi transitions the neutral edge at  $7.1 \text{ keV}$  could have an influence on the detection of lower velocity systems.

Of the sources included in the sample CBS 126, Mrk 766 (OBSIDs 701035010, 701035020), NGC 3227 (OBSIDs 703022010, 703022020, 703022030), NGC 4051 (OBSIDs 700004010, 703023010) and NGC 4151 all have at least one significantly detected absorption line at a rest-frame energy which is consistent at the 90 percent level with the neutral Fe K-shell edge from the reflection component. The absorption in Mrk 766, NGC 3227 and NGC 4051 is manifested by two lines with a common velocity shift equal to that expected for Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  and are therefore unlikely to be affected by the presence of an edge. However, the absorption in CBS 126 and NGC 4151, which only comprises a single detected profile, could possibly be affected. Even so, given that the iron K-shell edge structure is already self-consistently accounted for by the *reflectionx* model and in *xstar*, suggests that the residual profiles detected near the Fe K-shell edge in these sources is a real additional component rather than the residuals left by an inadequately fitted edge. Alternatively, the residuals could be due to a partially-covering absorber with a low outflow velocity. We investigate this possibility further in Appendix C. Finally, we note that while NGC 5506 shows evidence for a highly ionized Fe xxv emission line at  $\sim 6.63 \text{ keV}$  the absorption trough detected at  $\sim 9.2 \text{ keV}$  is not consistent with the Fe xxv K-shell edge and the two features are unlikely to be directly associated.

In addition, there are a few other complications which could have an effect on the identification of blue-shifted absorption profiles at  $E > 7 \text{ keV}$ . In particular, because different combinations of *xstar*

**Table 3.** *xstar* parameters for Fe K absorbers. Notes: (1) source name; (2) *Suzaku* observation ID; (3) logarithm of absorber column density, in units of  $\text{cm}^{-2}$ ; (4) logarithm of the ionization parameter, in units of  $\text{erg cm s}^{-1}$ ; (5) measured absorber outflow velocity, in units of  $v/c$ . Negative outflow velocities indicate a net red-shift; (6) change in  $\chi^2/\nu$  when the absorber removed from the best-fitting model; (7) corresponding absorber significance according to the *F*-test.

Source (1)	OBSID (2)	$\log N_{\text{H}}$ (3)	$\log \xi$ (4)	$v_{\text{out}}$ (5)	$\Delta\chi/\Delta\nu$ (6)	$P_{\text{F}}$ (7)
3C 111	703034010	$23.18^{+0.19}_{-0.24}$	$4.63^{+0.22}_{-0.18}$	$0.039 \pm 0.005$	18.7/3	99.97
	705040010	$22.43^{+0.20}_{-0.31}$	$4.13^{+0.27}_{-0.14}$	$0.105 \pm 0.008$	11.6/3	99.08
	(mean)	$22.95^{+0.28}_{-0.39}$	$4.45^{+0.35}_{-0.23}$	$0.072 \pm 0.005$	–	–
3C 390.3 <sup>Δ</sup>	701060010	>23.68	>5.46	$0.145 \pm 0.007$	11.7/2	>99.11
4C+74.26 <sup>Δ</sup>	702057010	>21.79	$4.06^{+0.45}_{-0.45}$	$0.185 \pm 0.026$	14.6/3	99.76
APM 08279+5255	stacked[all]	$23.00^{+0.07}_{-0.16}$	$3.41^{+0.08}_{-0.08}$	$0.139 \pm 0.012$	27.3/3	>99.99
		23.00*	$3.61^{+0.23}_{-0.14}$	$0.431 \pm 0.019$	14.3/2	99.59
	(mean)	$23.00^{+0.12}_{-0.12}$	$4.13^{+0.23}_{-0.17}$	$0.285 \pm 0.011$	–	–
CBS 126	705042010	>23.73	$4.77^{+0.26}_{-0.17}$	$0.012 \pm 0.006$	28.9/3	>99.99
ESO 103-G035 <sup>Δ</sup>	703031010	>21.90	$4.36^{+1.19}_{-1.19}$	$0.056 \pm 0.025$	33.2/3	>99.99
MCG -6-30-15	stacked[all]	$22.16^{+0.08}_{-0.08}$	$3.64^{+0.05}_{-0.06}$	$0.007 \pm 0.002$	103.0/3	>99.99
MR 2251–178	705041010	$21.54^{+0.22}_{-0.20}$	$3.26^{+0.12}_{-0.12}$	$0.137 \pm 0.008$	42.4/3	>99.99
Mrk 279	704031010	$23.38^{+0.25}_{-0.31}$	$4.42^{+0.15}_{-0.27}$	$0.220 \pm 0.006$	14.5/3	99.77
Mrk 766	701035010	$22.64^{+0.20}_{-0.27}$	$4.02^{+0.21}_{-0.13}$	$0.061 \pm 0.008$	14.5/3	99.71
	701035020	$22.76^{+0.10}_{-0.13}$	$3.67^{+0.06}_{-0.06}$	$0.017 \pm 0.004$	63.2/3	>99.99
	(mean)	$22.70^{+0.15}_{-0.19}$	$3.86^{+0.13}_{-0.10}$	$0.039 \pm 0.006$	–	–
NGC 1365	702047010	$23.92^{+0.03}_{-0.03}$	$3.88^{+0.06}_{-0.07}$	$0.014 \pm 0.001$	4125.1/3	>99.99
	705031010	$23.40^{+0.10}_{-0.12}$	3.88*	$0.002 \pm 0.002$	191.8/3	>99.99
	(mean)	$23.73^{+0.06}_{-0.06}$	$3.88^{+0.06}_{-0.07}$	$0.008 \pm 0.001$	–	–
NGC 3227	703022010	$22.74^{+0.11}_{-0.13}$	$3.89^{+0.08}_{-0.11}$	<0.002	63.1/3	>99.99
	703022030	$22.59^{+0.16}_{-0.20}$	$3.86^{+0.13}_{-0.14}$	$0.007 \pm 0.004$	21.0/3	>99.99
	703022050	$22.62^{+0.12}_{-0.14}$	$3.89^{+0.11}_{-0.12}$	$0.011 \pm 0.004$	37.9/3	>99.99
	(mean)	$22.66^{+0.13}_{-0.15}$	$3.88^{+0.19}_{-0.21}$	$0.005 \pm 0.004$	–	–
NGC 3516	100031010	$22.56^{+0.14}_{-0.17}$	$3.84^{+0.11}_{-0.10}$	$0.004 \pm 0.002$	141.9/3	>99.99
NGC 3783	701033010	$21.75^{+0.15}_{-0.20}$	$3.50^{+0.13}_{-0.08}$	<0.005	16.2/3	99.90
	704063010	$21.83^{+0.09}_{-0.10}$	$3.45^{+0.06}_{-0.04}$	<0.008	48.2/3	>99.99
	(mean)	$21.79^{+0.16}_{-0.23}$	$3.48^{+0.14}_{-0.90}$	<0.007	–	–
NGC 4051	700004010	$22.78^{+0.08}_{-0.09}$	$4.05^{+0.05}_{-0.05}$	$0.020 \pm 0.002$	80.3/3	>99.99
	703023010	$22.80^{+0.11}_{-0.13}$	$4.94^{+0.09}_{-0.09}$	$0.015 \pm 0.002$	20.8/3	99.99
	(mean)	$22.79^{+0.10}_{-0.11}$	$4.00^{+0.07}_{-0.07}$	$0.018 \pm 0.001$	–	–
NGC 4151 <sup>Δ</sup>	701034010	>21.74	$3.69^{+0.64}_{-0.64}$	$0.055 \pm 0.023$	54.8/3	>99.99
NGC 4395	702001010	$22.84^{+0.21}_{-0.25}$	$3.92^{+0.16}_{-0.15}$	<0.001	19.1/3	99.97
NGC 5506	stacked[all]	$23.22^{+0.19}_{-0.28}$	$5.04^{+0.29}_{-0.17}$	$0.246 \pm 0.006$	16.8/3	99.94
PDS 456	701056010	$23.04^{+0.08}_{-0.08}$	$4.19^{+0.15}_{-0.14}$	$0.253 \pm 0.008$	20.4/3	>99.99
		23.04*	4.19*	$0.292 \pm 0.009$	12.8/1	99.68
	705041010	23.04*	$3.93^{+0.02}_{-0.02}$	0.253*	20.4/3	>99.99
		23.04*	3.93*	0.292*	11.6/1	99.85
	(mean)	$23.04^{+0.08}_{-0.08}$	$4.06^{+0.15}_{-0.14}$	$0.273 \pm 0.006$	–	–
SW J2127.4+5654	702122010	$22.78^{+0.23}_{-0.34}$	$4.16^{+0.29}_{-0.13}$	$0.231 \pm 0.006$	11.5/3	99.24

\* indicates that a parameter was tied during spectral fitting.

<sup>Δ</sup> denotes sources with *xstar* solutions which are degenerate at the 90 per cent level (i.e. the  $\chi^2$  statistic for the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  solutions differs by  $\Delta\chi^2 \leq 2.71$ ). In these cases the reported values are averaged over the solutions and the errors are inferred as half the range between the absolute maximum and minimum values. Significances reported in columns (6) and (7) refer to the least significant of the valid solutions. See the text for further details.

parameters (i.e.  $N_{\text{H}}$ ,  $\log \xi$ ) can yield equivalent solutions to the  $\chi^2$  distribution at differing red-shifts there can be a level of degeneracy when identifying the ion responsible for a discrete absorption trough. This effect is not particularly significant in sources where

there are two absorption lines at a common velocity shift because the absorber is determined by the joint constraint of fitting both profiles, but it can become important in the instances where the absorption is manifested through a single trough. In these cases it can

be difficult to determine whether an absorption line is due to, for example, Fe  $\text{xxv}$  He $\alpha$  or Fe  $\text{xxvi}$  Ly $\alpha$ , which can therefore influence the inferred velocity of the outflow and hence any inferred absorber energetics.

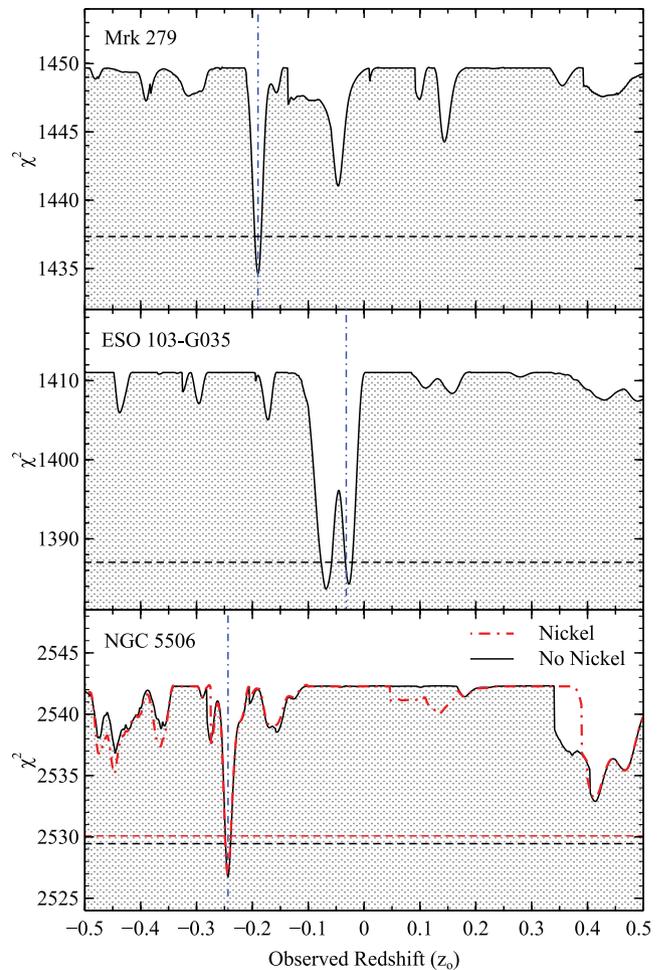
There are 10 sources (11 observations) in the sample in which a single absorption trough is detected at  $E > 7$  keV at the  $P_{\text{MC}} \geq 95$  per cent significance (see Table 2). In each of these observations we conducted a search for alternative *xstar* solutions by stepping the red-shift of the highly ionized *xstar* table through the Fe K band of each spectrum. This enabled the  $\chi^2$  minima for different *xstar* solutions at different velocity shifts, for various combinations of column density and ionization parameter, to be mapped. This process is analogous to that used by T11 and is useful when it comes to ascertaining whether there were any alternative fits to the absorption lines, and to check for degeneracies between line identifications.

For each source where only a single line is detected we took the best-fitting continuum model (including any necessary soft-band absorbers) and froze all model parameters bar those of the highly ionized absorption table and the normalization of the primary power-law. To map both blue- and red-shifted  $\chi^2$  minima the highly ionized *xstar* table was then stepped in  $\Delta z_0 = 10^{-3}$  increments within  $-0.5 \leq z_0 < 0.5$ , and contour plots were produced after each run. Example contour plots for Mrk 279 and ESO 103-G035, which illustrate the cases of non-degenerate and degenerate *xstar* solutions, respectively, are shown in Fig. 6. In Mrk 279 (Fig. 6, top panel) the stepping process yields a single valid *xstar* solution at a velocity shift corresponding to Fe  $\text{xxvi}$  Ly $\alpha$ . However, in ESO 103-G035 (Fig. 6, middle panel) there are two degenerate  $\chi^2$  minima found which are statistically equivalent at the 90 per cent level (i.e.  $\Delta\chi^2 \leq 2.71$ ). The measured parameters imply that the lowest velocity solution is associated with Fe  $\text{xxvi}$  K $\alpha$ , while the higher velocity solution is likely due to Fe  $\text{xxv}$  He $\alpha$  with some contamination from lower ionized species of iron. In this scenario it is difficult to unambiguously identify the responsible Fe ion and therefore gauge the appropriate outflow velocity.

Of the 10 sources in which the Fe K absorber is manifested by a single trough only 4 show evidence for degenerate *xstar* solutions (3C 390.3, 4C+74.26, ESO 103-G035 and NGC 4151). For these cases we report the mean *xstar* parameters as measured from all valid solutions in Table 3 to account for the uncertainty in line identifications, with the associated errors being taken as half of the range between the absolute and minimum values. The remaining 7 sources have only a single  $\chi^2$  minimum which corresponds to the line identification reported in Table 2. In sources with degenerate *xstar* solutions we report the most conservative line identification (i.e. associated with Fe  $\text{xxvi}$  Ly $\alpha$ ) in Table 2.

#### 4.6.2 The influence of nickel

In principle, the presence of nickel could also complicate our line identifications above 7.1 keV. In particular, the K $\alpha$  lines of Ni  $\text{xxvii}$  and Ni  $\text{xxviii}$ , which are expected at rest-frame energies of  $E \sim 7.78$  keV and  $E \sim 8.09$  keV, respectively, could offer a more energetically conservative identification for the highest energy absorption lines by virtue of requiring a lower blue-shifted velocity. A total of six sources have absorption lines detected at energies in the regime where nickel could complicate line identifications (Table 2). Because of uncertainties in the atomic rates of nickel<sup>13</sup>



**Figure 6.** Example plots showing the  $\chi^2$  statistic versus observed redshift of the *xstar* absorber ( $z_0$ ). Plots are shown for Mrk 279 (top panel), ESO 103-G035 (middle panel) and NGC 5506 (bottom panel). In all panels the black solid line shows the  $\chi^2$  confidence contour and underlying  $\chi^2$  distribution. In the bottom panel the black and red solid lines show the confidence contours for *xstar* grids both with (dot-dashed red) and without (solid black) nickel included at solar abundances. In all panels the horizontal dashed lines show the 90 per cent confidence level for one interesting parameter, while the vertical dash-dotted line shows the best-fitting absorber red-shift.

its abundance is set to zero in *xstar* by default which means that any additional solutions to the  $\chi^2$  distribution due to it would not be picked up during the red-shift stepping. In order to robustly search for plausible *xstar* solutions which may be associated with nickel we therefore regenerated the *xstar* tables using the same assumed parameters as before (see Section 4.6), but this time with nickel included at solar abundances, and again searched for alternative *xstar* solutions using the method outlined previously.

The bottom panel of Fig. 6 shows a comparison of the contour plots obtained for NGC 5506 when nickel is set to zero or solar in the *xstar* table. As expected given the  $Z_{\text{Ni}}/Z_{\text{Fe}} \sim 0.05$  abundance ratio at solar abundances (i.e. Grevesse & Sauval 1998), in terms of valid *xstar* solutions for NGC 5506 there are no tangible differences between the two grids with Fe  $\text{xxvi}$  Ly $\alpha$  being the only valid identification in both cases. Similar is also true for the other five sources with absorption lines detected at  $E \gtrsim 7.78$  keV, with no valid solutions corresponding to nickel being found in any of the cases. Moreover, it should be noted that in order to achieve the measured

<sup>13</sup> <http://heasarc.gsfc.nasa.gov/xstar/docs/html/node74.html>

**Table 4.** Fraction of detected outflows.

Line(s)	Sources (spectra)
Fe xxv He $\alpha$	2 (3)
Fe xxvi Ly $\alpha$	9 (10)
Fe xxv–xxvi	7 (12)
Multi $v_{\text{out}}$	2 (3)
Total	20 (28)

$EW$  of the absorption lines the required column density would be unreasonably large, i.e.  $\log(N_{\text{H}}/\text{cm}^{-2}) > 24$ , which would absorb the observed continuum far beyond that which is observed in the spectra.

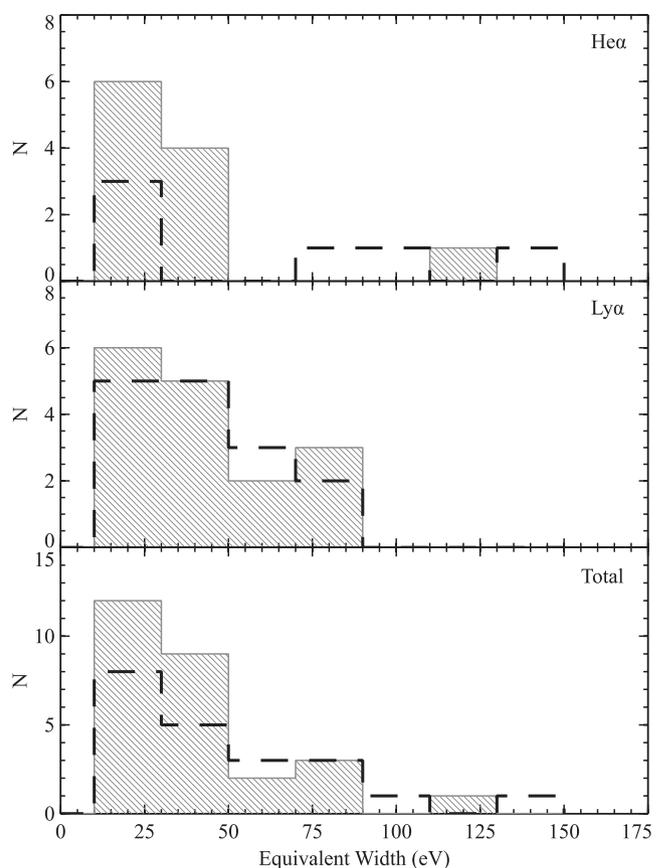
Therefore, bar the extraordinary case where it is  $\gtrsim 20$  times over-abundant relative to solar values, Ni has a negligible influence on the Fe K band and identifying the observed absorption lines with K-shell resonance lines of Fe is secure. We report the best-fitting `xstar` parameters, including the measured outflow velocities relative to the host galaxy, in Table 3. The statistical significances as per the  $\Delta\chi^2$  and the  $F$ -test are also reported; in all cases the addition of an `xstar` grid to fit the observed lines improves the fit by at least the 99 per cent level by the  $F$ -test. Note that since some sources either have more than one absorption complex detected in a single spectrum (i.e. APM 08279+5255 and PDS 456), or have absorption detected in more than one epoch (i.e. Mrk 766), we also report the mean absorption parameters on a per-source basis. The mean values are used in all subsequent analysis to prevent sources with multiple Fe K absorber detection from over-weighting the resultant parameter distributions which are discussed in Section 5.2.

## 5 RESULTS

### 5.1 Line detection rate and phenomenology

A total of 20/51 sources (in 28/73 fitted spectra, corresponding to  $\sim 40$  per cent of the total sample) show evidence for highly ionized absorption lines in their *Suzaku* spectra at a Monte Carlo significance of  $P_{\text{MC}} \geq 95$  per cent, with 18/20 of these outflows also robustly detected at  $P_{\text{MC}} \geq 99$  per cent significance (see Table 2). The tally of line phenomenologies is shown in Table 4. Of the 28 observations with individually detected Fe K absorbers there are 10 which are consistent with having Fe xxvi Ly $\alpha$  as the dominant Fe ion, 3 where Fe xxv He $\alpha$  (and/or lower ionization species of iron) is the main contributor, 12 absorbers with both Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  lines with a common outflow velocity and a further 3 that have two absorption components with different outflow velocities. Taking into account that some sources have absorption lines detected in more than one observation, on a per-source basis this corresponds to 2/20 having Fe xxv He $\alpha$  absorption, 9 with Fe xxvi Ly $\alpha$  absorption<sup>14</sup>, 7 with both Fe xxv and Fe xxvi Ly $\alpha$  absorption and 2 having multi- $v_{\text{out}}$  systems (out of 20). Interestingly, 8/9 sources with just a single Fe xxvi Ly $\alpha$  absorption line have outflow velocities which exceed the  $v_{\text{out}} \sim 10\,000$  km s<sup>-1</sup> threshold employed by T10A when identifying ultra-fast outflow (UFO) systems, while only 1 of the 6 sources with Fe xxv–xxvi absorption

<sup>14</sup> In reality the number of Fe xxvi Ly $\alpha$  systems could be somewhat lower than this value due to the four AGN where the absorption is equally well fitted by Fe xxv He $\alpha$  at a slightly lower velocity (i.e. see Section 4.6.1).



**Figure 7.** Histogram showing the mean  $EW$ s of detected highly ionized absorption lines. The top and middle panels show the distributions for Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  transitions, respectively, while the bottom panel shows the total distribution of all detected  $1s \rightarrow 2p$  absorption lines including those which are detected as a blend of the two K $\alpha$  transitions. For sources with more than one line detected in multiple observations the mean  $EW$ s have been used. The grey diagonally shaded area and the dashed black line show the distributions found in this work and with *XMM-Newton*, respectively. The *XMM-Newton* distributions are independently created using the values listed in table A.3 in T10A.

exhibits a mean velocity which exceeds this threshold value. This is consistent with the view that higher ionized outflows originate closer to the central AGN and therefore have a higher outflow velocity, and is a point which will be further discussed in a companion paper (Gofford et al., in preparation).

Histograms of  $EW$  for the absorption lines are shown in the top and middle panels of Fig. 7, respectively. Note that in these plots the individual line profiles in a Fe xxv–xxvi pair are considered separately for the sake of clarity, and only those lines which have been individually detected at a Monte Carlo significance of  $P_{\text{MC}} \geq 95$  per cent are considered. The histogram in the bottom panel shows the  $EW$  distribution for *all* statistically significant absorption lines, including those which are consistent with being a blend of the Fe xxv–xxvi K $\alpha$  transitions. Measured  $EW$ s span from a few tens of eV up to  $\sim 130$ – $140$  eV (Fig. 7, bottom panel), in a distribution which is consistent with the curve of growth analysis conducted by T11. The mean  $EW$ s of the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  absorption lines are  $\sim 36$  and  $\sim 38$  eV, respectively, with the total mean  $EW$  for all detected profiles (i.e. included blended ones)  $\sim 40$  eV. Importantly, the observed distribution of absorption line  $EW$ s is broadly consistent with that found by T10A using *XMM-Newton* (marked with a dashed black line in Fig. 7).

As noted by T10A a useful quantity to calculate is the global probability that all of the observed lines are due to purely statistical shot noise, which can be done using the binomial distribution. For an event with null probability  $p$  the chance probability of the event  $n$  happening in  $N$  trials is given by

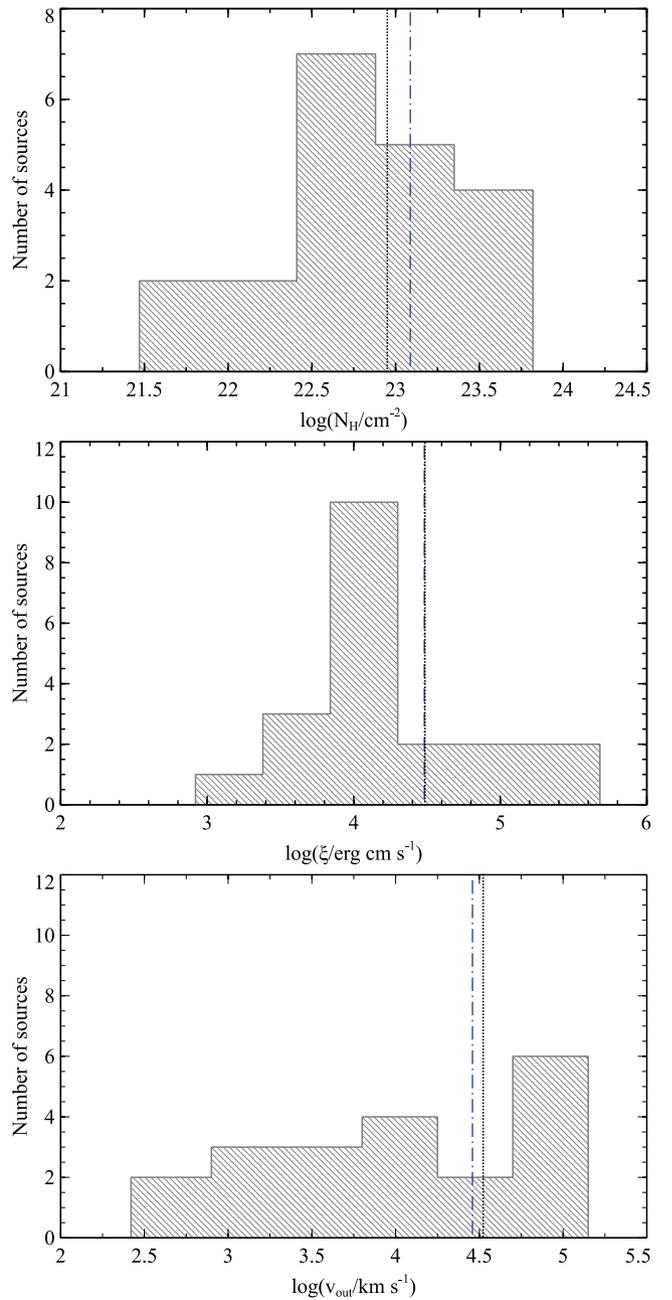
$$P(n; N, p) = \frac{N!}{n!(N-n)!} p^n (1-p)^{N-n}. \quad (1)$$

For  $n = 20$  Fe K-band absorption line systems detected in  $N = 51$  sources at a Monte Carlo significance of  $P_{MC} \geq 95$  per cent the probability of one of these absorption systems being due to shot noise can be taken as  $p < 0.05$ . On a per-source basis the probability of all of the observed absorption systems being associated with noise is then very low, with  $P < 2 \times 10^{-13}$ , which further reduces to  $P < 5 \times 10^{-18}$  when considering the fact that some sources have lines detected in more than one epoch. This suggests that the observed lines are very unlikely to be associated with simple statistical fluctuations in the spectra. Moreover, it is important to remember that because the vast majority of absorption lines are detected at  $P_{MC} > 95$  per cent significance these probabilities only represent conservative lower limits on the global probability of all lines being false detections.

## 5.2 Absorber properties

The *xstar* parameter distribution as measured with *Suzaku*, again plotted using the mean absorber parameters averaged over all observations, is shown in Fig. 8. Overall, the general distributions and mean parameter values are broadly consistent with those found by T10A using *XMM-Newton*. The Fe K absorbers detected with *Suzaku* cover a wide range of column densities,  $21.5 < \log(N_{\text{H}}/\text{cm}^{-2}) \leq 24.0$ , with a peak in the distribution at  $\log(N_{\text{H}}/\text{cm}^{-2}) \approx 22\text{--}23$ . As shown by the dot-dashed (blue) and dotted (black) lines, which show the mean  $\log(N_{\text{H}}/\text{cm}^{-2})$  value as found with *Suzaku* and *XMM-Newton*, respectively, the mean is  $\log(N_{\text{H}}/\text{cm}^{-2}) \approx 23$  for both samples. From the middle panel the ionization parameters are in the interval  $2.5 < \log(\xi/\text{erg cm s}^{-1}) \leq 6$  with the significant fraction of Fe  $\text{xxv}\text{--}\text{xxvi}$  pair systems, which persist over only a relatively narrow range in ionization parameter (see the curve of growth analysis by T11), leading to a peak in the distribution at  $\log(\xi/\text{erg cm s}^{-1}) \approx 4$ . Again, as shown by the vertical lines the mean ionization parameter in both samples is almost identical at  $\log(\xi/\text{erg cm s}^{-1}) \approx 4.5$ . The detection of relatively low-ionization material in the Fe K band, i.e.  $\log(\xi/\text{erg cm s}^{-1}) \sim 2.5\text{--}3.0$ , is particularly interesting and suggests that high-velocity absorption could feasibly be detected at softer X-ray energies through weak, moderately ionized iron lines. Moreover, the detection of a small fraction of absorbers with  $\log(\xi/\text{erg cm s}^{-1}) \geq 5$  in both this work and in the T10A sample raises the possibility that material may be present in some sources which is so highly ionized that even iron is not detectable through spectroscopy. If this is the case then the fraction of sources with Fe K absorption ( $\sim 40$  per cent) may represent a lower limit on the number of sources with intrinsic nuclear outflows along the line of sight.

The  $\log(v_{\text{out}}/\text{km s}^{-1})$  distribution (bottom panel) appears to be relatively continuous over a broad range of velocities, ranging from as low as  $v_{\text{out}} < 1500 \text{ km s}^{-1}$  up to  $v_{\text{out}} \approx 100\,000 \text{ km s}^{-1}$ . The median outflow velocity is  $v_{\text{out}} \sim 17\,000 \text{ km s}^{-1}$  ( $\sim 0.056 c$ ). Ninety per cent of the detected outflows have  $v_{\text{out}} \geq 1500 \text{ km s}^{-1}$  which makes the absorption detected at Fe K almost systematically faster than the traditional soft X-ray warm absorber. Only



**Figure 8.** Histogram showing the overall distributions of mean absorber parameters for each source: (a) logarithm of the mean column density; (b) logarithm of the mean ionization parameter; (c) logarithm of the mean outflow velocity. The dot-dashed (blue) and dotted (black) vertical lines indicate the mean value of the *Suzaku* and *XMM-Newton* analyses, respectively.

NGC 3227, NGC 4395 and NGC 3783 have Fe K outflows which are consistent with having no outflow velocity. From Table 5 the distributions of outflow velocity appear to be very similar between both the *Suzaku* and *XMM-Newton* samples, with the bulk of outflows in both samples having  $v_{\text{out}} > 10\,000 \text{ km s}^{-1}$ . The *Suzaku* sample does appear to have slightly more low-intermediate velocity systems but, given the low number statistics involved, the differences are probably not significant. An interesting possibility is that both *Suzaku* and *XMM-Newton* may be subject to an instrumental bias against the detection of low-velocity absorption systems. At CCD resolution the presence of Fe  $\text{xxv}$  and/or Fe  $\text{xxvi}$  emission lines could

**Table 5.** Outflow velocity comparison.

Velocity (km s <sup>-1</sup> )	<i>Suzaku</i>	<i>XMM-Newton</i>
No outflow	2/20	2/19
$0 < v_{\text{out}} \leq 10\,000$	6/20	2/19
$v_{\text{out}} > 10\,000$	12/20	15/19
$v_{\text{out}} \geq 30\,000$	8/20	9/19

mask the signature of low-velocity absorption systems and thereby introduce a selection effect against their detection. The only way to reveal the presence of such absorption systems would be through future observations with, for example, the calorimeter aboard *Astro-H*, which would have sufficient resolution to distinguish between individual emission and absorption components at low velocity.

Even so, the similarity between all of the measured *Suzaku* and *XMM-Newton* parameter distributions can be quantitatively assessed using the Kolmogorov–Smirnov two-sample test (K-S test) which uses the maximum differences between the cumulative fractional distribution of two sets of data to determine the probability that they are both drawn from the same parent sample. In this case such a test can be used to quantify the level at which the column density, ionization parameter and outflow velocity distributions as measured with both *Suzaku* and *XMM-Newton* are in agreement. For a null hypothesis that each of the *Suzaku* and *XMM-Newton* distributions is drawn from the same parent sample we are unable to conclusively rule out the null hypothesis in any of the three cases at greater than the 90 per cent confidence level.

## 6 DISCUSSION

### 6.1 Detailed sample comparison

#### 6.1.1 Radio-quiet sources

In this work we have detected Fe K outflows in 17 radio-quiet AGN and 3 in radio-loud. Of the detections in radio-quiet sources, 9/17 also have had observations included in the *XMM-Newton* outflow sample, with 6/9 of these cases having Fe K outflows confirmed by the T10A analysis (i.e. Mrk 766, NGC 4051, NGC 3516, NGC 3783, NGC 4151, Mrk 279). In the remaining 3 sources, namely NGC 5506, MCG -6-30-15 and NGC 3227, the *Suzaku* outflow detections are not confirmed by T10A. Even so, given that in the case of NGC 3227 the Fe K absorption lines are detected at  $P_{\text{MC}} > 99.9$  per cent significance in three separate *Suzaku* observations, and in MCG -6-30-15 the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  absorption lines present are both statistically resolved at  $P_{\text{MC}} > 99.9$  per cent confidence, the *Suzaku* line detections alone still imply very robust absorber detections in both of these sources. A further three sources with absorption detected in their *Suzaku* spectra have had statistically significant absorption detected in *XMM-Newton* observations which were not included in the T10A outflow sample (NGC 1365, Risaliti et al. 2005; PDS 456, Reeves et al. 2003, 2009; APM 08279+5255, Chartas et al. 2009). Furthermore, weak hard-band absorption has been noted in the *XMM-Newton* spectra of MCG -6-30-15 by several authors (e.g. Fabian et al. 2002; Vaughan & Fabian 2004; Nandra et al. 2007; Miller et al. 2008). In particular, Nandra et al. (2007) reported the presence of moderately significant (i.e. just below 99 per cent confidence) Fe xxv He $\alpha$  absorption in XMM\_OBSID: 00297401010, but this detection was not significantly replicated by T10A and highlights the need for de-

tailed broad-band spectral models when assessing for the presence of highly ionized absorption lines.

Therefore a total of 9/17 (~53 per cent) of the sources with outflows detected by *Suzaku* have also previously been detected in *XMM-Newton* data. In addition to these, the absorbers in MR 2251–178 and MCG -6-30-15 are also corroborated on the basis of their *Chandra* High Energy Transmission Grating (HETG) data. Gibson et al. (2005) detected a resolved Fe xxv He $\alpha$  absorption line at  $E \sim 7.25$  keV in the HETG spectrum MR 2251–178 ( $v_{\text{out}} \sim 13\,000$  km s<sup>-1</sup>) while Young et al. (2005) found variable absorption lines due to Fe xxv and Fe xxvi Ly $\alpha$  in MCG 6-30-15 with  $v_{\text{out}} \sim 2000$  km s<sup>-1</sup>. In both cases the outflow velocities are consistent with those found in our *Suzaku* analysis, and the detection in the HETG spectrum of MCG -6-30-15 overcomes the ambiguities which remained on the basis of the *XMM-Newton* data and independently corroborates the *Suzaku* line detections. Fe K absorption has also been reported in the *Chandra* spectrum of APM 08279+5255 (Chartas et al. 2002) making it the only source in the sample to have its Fe K outflow independently detected in all three observatories at a high significance level.

Thus, a total of 11/17 (~65 per cent) sources with Fe K outflows detected with *Suzaku* have also had outflows reported elsewhere in the literature with either *XMM-Newton* or *Chandra*. Of the remaining 6 sources, SWIFT J2127.4+5654 has been observed by *XMM-Newton* twice: once in 2009 (~24 ks) and once in 2010 (~131 ks), although neither observation has been published at the time of writing. CBS 126 have not been observed by either *XMM-Newton* or *Chandra*, while ESO 103-G035 only has a 13 ks *XMM-Newton* observation available which is insufficient to test for Fe K absorption. This leaves only three AGN, NGC 5506, NGC 3227 and NGC 4395, with sufficient observations available which do not have independent confirmations with other observatories; however, as discussed previously, NGC 3227 does have multiple detections with *Suzaku* which suggests a robust detection.

#### 6.1.2 Radio-loud sources

There are six BLRGs in our *Suzaku* sample. This includes all of the sources which were part of the T10B radio-loud UFO case study (i.e. 3C 111, 3C 120, 3C 445, 3C 390.3 and 3C 382), as well as 4C+74.26 which was not part of their analysis. T10B state that 3/5 of the BLRGs in their sample, namely 3C 111, 3C 120 and 3C 390.3, show evidence for highly ionized absorption lines in their *Suzaku* spectra, with that found in 3C 111 also being confirmed in more recent subsequent observations (Tombesi et al. 2011b).

We are able to confirm the absorption lines detected in both 3C 111 and 3C 390.3 at a similar confidence level. However, while we are unable to directly confirm the line detection reported in 3C 120 we note that by adding two narrow ( $\sigma = 10$  eV) Gaussians at the rest-frame energies for the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  lines reported by T10B (7.25 and 7.54 keV, respectively) we can place lower limits of  $EW_{\text{He}\alpha} > -8$  eV and  $EW_{\text{Ly}\alpha} > -9$  eV for the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  absorption lines, respectively. These limits are consistent with those reported by T10B at the 90 per cent level. Furthermore, we have also detected a high-velocity outflow ( $v_{\text{out}} \sim 0.22c$ ) in the BLRG 4C+72.26 and note that an Fe K outflow is also detected in the *Chandra* Low Energy Transmission Grating (LETG) data for 3C 445 (Reeves et al. 2010; see also Braito et al. 2011). Thus the number of BLRGs with high-velocity outflows could tentatively rise to 4/6 if observations from multiple observatories are included, suggesting that such outflows could be an important component in

radio-loud sources. It is important to note, however, that the current sample size of seven *Suzaku* observed radio-loud AGN is insufficient for a statistical study into the prevalence of Fe K absorption in radio-loud sources. That being said, Tombesi et al. (in preparation) are currently performing a detailed analysis on a large sample of  $\sim 30$  radio-loud AGN observed using *XMM-Newton* and *Suzaku* which will robustly assess for the prevalence of Fe K absorption in these objects.

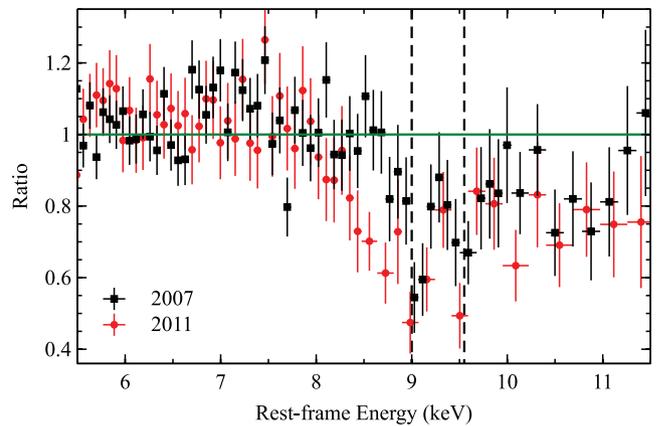
### 6.1.3 Recent *Suzaku* samples

The analysis of another large sample of *Suzaku* selected AGN has been published by Patrick et al. (2012). Whilst the primary aim of their study was to assess the properties of Fe K emission lines and black hole spin, the authors also report the detection of highly ionized Fe K absorption in 14/46 ( $\sim 30$  per cent) of their objects. The two samples have 43 AGN in common. This includes 15 AGN in which we have detected Fe K absorption, of which 11 are confirmed by the Patrick et al. (2012) analysis. The four objects where our outflow detections are not corroborated are 3C 390.3, 4C+74.26, NGC 5506 and SWIFT J2127.4+5654. The outflows that we detect in three of these objects (3C 390.3, 4C+74.26 and SWIFT J2127.4+5654) are only marginal detections from Monte Carlo simulations, i.e.  $95 \text{ per cent} \leq P_{\text{MC}} < 99 \text{ per cent}$ , which suggests that the lines may not be visually apparent in the raw data without a statistically driven analysis. The outflow in NGC 5506 is detected with  $P_{\text{MC}} = 99.8 \text{ per cent}$  which suggests that the discrepancy is likely a result of the method in which the data sets were analysed. Indeed, our detection of Fe xxvi Ly $\alpha$  in this source is taken from the time-averaged (stacked) *Suzaku* spectrum while Patrick et al. (2012) analysed the observations separately. The outflow in NGC 5506 may be intrinsically weak in the individual *Suzaku* epochs, and only become apparent when they are time-averaged due to higher S/N.

An interesting outcome of the Patrick et al. (2012) analysis is that they report outflow detections in three sources (3C 445, NGC 5548 and PG 1211+143) which we do not find in this work. Outflows in 3C 445 and PG 1211+143 have been detected with other observatories (3C 445: Reeves et al. 2010, PG 1211+143: Pounds et al. 2003b) and they are known to be intrinsically variable, being weakest in their respective *Suzaku* observations (Reeves et al. 2008b; Braito et al. 2011). In NGC 5548 we note that while there is a weak absorption trough detected at the energy expected for Fe xxv He $\alpha$  (see appropriate entry in Appendix B) it falls below our detection criteria. Similar is also true for the *Suzaku* line detection in PG 1211+143, which is only detected at 90 per cent significance. However, when compared to the sample of Patrick et al. (2012) the broader outcomes are entirely consistent. Both analyses find that there is a strong peak in the ionization parameter distribution at  $\log(\xi/\text{erg cm s}^{-1}) \sim 4\text{--}4.5$ , with the associated velocities spanning a continuous range  $\sim 0 \lesssim v_{\text{out}} \lesssim 100\,000 \text{ km s}^{-1}$ . Therefore, the two samples are in complete agreement when it comes to the overall properties of the Fe K absorbers.

## 6.2 Evidence for complex variability and absorber structure?

In addition to broad-band spectral variability there is also compelling evidence for intricate variability in the Fe K absorber itself. In a number of cases this variability can be simultaneously and self-consistently fitted using *xstar* absorption tables which strongly implies an atomic origin rather than an association with spectral



**Figure 9.** Ratio plot comparison of the 2007 (black squares) and 2011 (red circles) *Suzaku* observations of PDS 456 in the Fe K band. The dashed vertical lines indicate the position of the two blended Fe xxv–xxvi absorption troughs present in both spectra, at rest-frame energies of  $\sim 9.0$  and  $\sim 9.6$  keV, respectively. There is a noticeable broadening in the left-hand absorption profile in the 2011 observation which is likely caused by a drop in absorber ionization state, and an associated increase in the Fe xxv He $\alpha$  contribution to the absorption profile. See the text and Reeves et al. (in preparation) for further details.

fluctuations and photon statistics. In NGC 3227 transient Fe K absorption lines are detected in three of the six *Suzaku* epochs, with the line parameters varying on a  $\Delta t \sim 2$  week time-scale. Similar is also true in Mrk 766, PDS 456 and NGC 1365, which all exhibit absorption line variability albeit over a longer (i.e.  $\Delta t \sim \text{year}$ ) base-line. In both Mrk 766 and PDS 456 the absorber appears to change in ionization states between the *Suzaku* epochs, with the outflow in Mrk 766 also showing a decrease in outflow velocity. Here, we discuss both sources in greater detail.

Mrk 766 and PDS 456 have both long been known to harbour Fe K absorbers (Mrk 766: Pounds et al. 2003a; Turner et al. 2007; Miller et al. 2008; Risaliti et al. 2011. PDS 456: Reeves et al. 2000, 2003, 2009). The range of rest-frame energies for the absorption profiles measured with *Suzaku* in each source are consistent with those previously published in the literature. In Mrk 766 we find a common velocity shift of  $v_{\text{out}} \sim 6000\text{--}18\,000 \text{ km s}^{-1}$  for the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  lines, which is consistent with the  $v_{\text{out,XMM}} \sim 3000\text{--}16\,000 \text{ km s}^{-1}$  reported previously in *XMM-Newton* data (Miller et al. 2007; Turner et al. 2007; Risaliti et al. 2011). Similarly, from our simultaneous analysis of the 2007 and unpublished 2011 *Suzaku* observations of PDS 456 we find two absorption troughs which persist across both sequences (see Fig. 9). These troughs – which are likely due to blends of the Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  lines owing to their breadth – have a mean inferred  $\langle v_{\text{out}} \rangle \sim 0.3 c$ , which is consistent with that reported by Reeves et al. (2009) in their initial analysis of the 2007 observation.

An interesting property of the outflows in both of these sources is that the parameters of the hard-band absorber appear to be linked to both the incident source flux and the broad-band absorption characteristics of the X-ray spectrum. On the basis of detailed time- and flux-resolved spectroscopy of Mrk 766 using *XMM-Newton*, Turner et al. (2007) noted that the EW of the Fe xxvi Ly $\alpha$  absorption line could be correlated with the intrinsic source flux, with the correlation not easily accounted for by simply appealing to changes in the ionization state of a constant column density absorber in response to changes in the incident continuum. This provided evidence for

complex absorber structure in Mrk 766. Further evidence for Fe K absorber complexity was suggested by Risaliti et al. (2011) who, from a re-analysis of the same *XMM-Newton* data, found that the presence of the highly ionized absorption profiles appeared to be contingent on the source being eclipsed by a number of high column density, i.e.  $\log(N_{\text{H}}/\text{cm}^{-2}) \gtrsim 23$ , partial covering clouds with a low ionization state. This led the authors to propose a scenario where multiple ‘cometary’ absorption clouds with a large gradient in both column density and ionization state were moving across the line of sight, thus giving rise to the varying absorption characteristics present in different time-slices (e.g.  $v_{\text{out}}$ ).

Our *Suzaku* analyses yield further evidence for a complex and dynamic structure in hard-band absorbers. In Mrk 766 the *EW*s of the Fe XXV He $\alpha$  and Fe XXVI Ly $\alpha$  lines present in the *Suzaku* spectrum are statistically consistent at the 90 per cent level but, as shown in the appropriate entries in Appendix B, there is a visual shift in the dominant iron ion of the absorber, with Fe XXVI Ly $\alpha$  ( $EW_{\text{FeXXVI}} = -40 \pm 17$  eV) appearing strongest in OBSID 701035010 and Fe XXV He $\alpha$  ( $EW_{\text{FeXXV}} = -60 \pm 14$  eV) strongest in OBSID 701035020, which is possibly associated with the  $\Delta \log(\xi/\text{erg cm s}^{-1}) \sim 0.4$  change in ionization parameter between the two observations. Furthermore, and as mentioned previously, the decrease in ionization is also met with a decrease in measured outflow velocity and the onset of strong Fe XXV emission which could be associated with a cloud/clump of material moving outside of the sight line. A similar effect is also apparent in PDS 456 where, despite there being no measurable change in the relative velocity of the two blended Fe XXV–XXVI absorption complexes, the absorber ionization state decreases by  $\Delta \log(\xi/\text{erg cm s}^{-1}) \sim 0.26$  in the 2011 observation relative to 2007 which leads to a slight visual broadening of the absorption profiles due to an increased Fe XXV He $\alpha$  contribution in the blended absorption line profile. Interestingly, the decreased ionization state of the absorbers in both Mrk 766 and PDS 456 occurs when the covering fraction of the soft X-ray absorber is at its highest which implies a dynamic clumpy structure to the outflow. Moreover, as noted by both Turner et al. (2007) and Behar et al. (2010), respectively, the opposite is also true in both sources, i.e. when the sources are in their least absorbed state, the discrete highly ionized absorption lines are not significantly present in the X-ray spectra of either source. This is consistent with the conclusions reached by Risaliti et al. (2011) in the context of stratified, partially-covering, absorption clumps needing to be present along the line of sight, and close to the continuum source, for highly ionized absorption lines to be observed.

Additional evidence for this supposition is provided through consideration of our spectral models to other sources in the sample. It is intriguing to note that statistically acceptable (i.e.  $\chi^2_{\text{reduced}} \approx 1.0$ ) partial covering absorption models are obtained for a curiously large fraction of the sources with detected outflows (12/20;  $\sim 60$  per cent) which further suggests that there may be a link between complex absorption geometries and highly ionized absorption lines. One possibility is where the partially-covering components represent denser clumps of material in an inhomogeneous highly ionized wind from the accretion disc, similar to the case reported in Mrk 766 (Miller et al. 2007; Turner et al. 2007; Risaliti et al. 2011). Indeed, inhomogeneous winds with stratified or filamentary ionization and density structure are expected as a natural consequence of the accretion process and are ubiquitously seen in both hydrodynamic (e.g. Proga & Kallman 2004; Kurosawa & Proga 2009; Sim et al. 2010) and magneto-hydrodynamic (e.g. Ohsuga et al. 2009; Ohsuga & Mineshige 2011) simulations of accretion discs. It is therefore not particularly surprising to find a possible link between clumpy

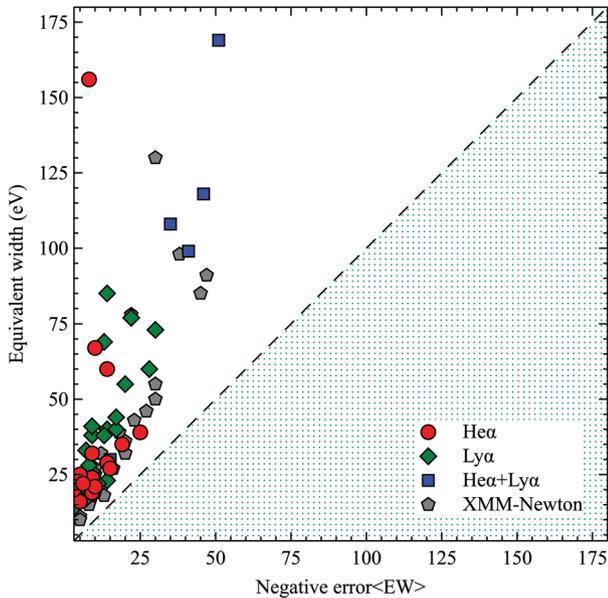
absorption dominated models and sources with high-velocity outflows.

### 6.3 On the claimed publication bias

The transient and variable nature of Fe K absorption lines has in part led to their true veracity to be questioned in the literature. In particular, Vaughan & Uttley (2008) suggested that there may be a publication bias at play in the reporting of both red- and blue-shifted features in the Fe K band, with only those observations with the strongest line detections being reported in the literature. Through a plot of ‘ $EW - \text{error}(EW)$ ’ Vaughan & Uttley (2008) showed that the  $EW/\text{error}(EW)$  ratio remained relatively constant (i.e. lines with larger  $EW$  have correspondingly larger 90 per cent errors) over a wide range of  $EW$ s in their sample of narrow lines in 38 sources collected from the literature. The authors also noted that since the vast majority of reported detections (prior to 2008, at least) were of relatively low significance (i.e. typically around  $2 - 3\sigma$ ), the lines may be spurious, and more consistent with merely being the strongest natural fluctuations in otherwise featureless spectra than real atomic features. Moreover, Vaughan & Uttley (2008) suggested that the conspicuous absence of any lines in the upper-left quadrant of the  $EW - \text{error}(EW)$  plot (which would correspond to stronger lines having smaller uncertainties) implied that the detection of the narrow velocity shifted lines was in some way inversely correlated to the statistical quality of the observation. Observations with longer exposures and better photon statistics only ever showed the weakest lines, which further enforced the possibility of a bias in the published observations.

However, these points were addressed through the systematic search conducted by T10A using *XMM-Newton*, at least in the context of blue-shifted Fe K absorption lines. By uniformly and comprehensively searching for such features in a complete sample of AGN observations, carefully reporting the fraction of detections to null detections, and assessing the statistical significance of any detected absorption lines through Monte Carlo simulations, T10A were able to robustly assess the fraction of AGN in the local universe that have Fe K outflows in a way which overcame any publication biases and accounted for the possibility of random fluctuations in the source spectra. Our work with the *Suzaku* outflow sample complements the findings of T10A, and lends further weight to the assertion that such outflows are an important intrinsic feature of the AGN X-ray spectrum in a large fraction of sources. In Fig. 10 we replicate the Vaughan & Uttley (2008)  $EW/\text{error}(EW)$  plot<sup>15</sup> used by T10A (see their fig. 8) to include data points corresponding to the absorption lines found in this work at a  $P_{\text{MC}} \geq 95$  per cent significance level. The figure concisely shows that the distribution of points from both the *Suzaku* and *XMM-Newton* samples diverges from that of the  $EW = \text{error}(EW)$  ‘detection line’, with both sets of data having a number of points which veer towards the important upper-left quadrant of the diagram indicating stronger lines with smaller uncertainties. Interestingly, the overall distribution of points for the *Suzaku* lines follows a similar trend to those obtained with *XMM-Newton* which implies that all points are drawn from the same parent population, as would be expected should we be

<sup>15</sup> Note that the axes in this plot use a linear scale rather than the logarithmic one as used by Vaughan & Uttley (2008). As noted by T10A the use of a logarithmic scale visually compresses the data points towards the  $EW = \text{error}(EW)$  threshold of non-detection.

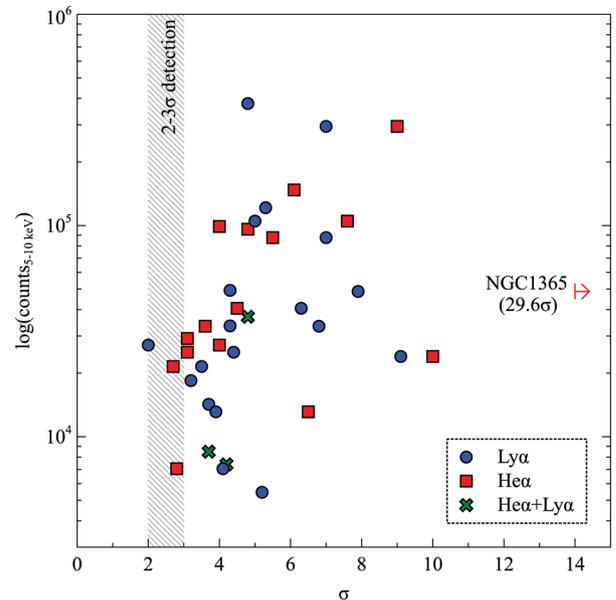


**Figure 10.** Plot of the measured  $EW$  of the blue-shifted absorption lines with respect to their negative 90 per cent error. The circles (red), diamonds (green) and squares (blue) correspond to the Fe xxv He $\alpha$ , Fe xxvi Ly $\alpha$  and blended Fe xxv–xxvi lines, respectively. Analogous values for the blue-shifted absorption lines found in the T10A *XMM–Newton* sample are shown as pentagons (grey). The diagonal dashed line shows the expected trend if  $EW = -\text{error}(EW)$  and the dotted (dark-cyan) area shows the area of non-detection (see the text for further details).

studying the same physical phenomenon which imprints real spectroscopic lines in the X-ray spectrum.

A further means of visually assessing for a publication bias, and for testing whether the absorption lines are consistent with random fluctuations or real features, is through a plot of line significance versus photon counts in the Fe K band (i.e. between 5 and 10 keV). Such a plot is shown in Fig. 11. For simplicity we grade line significance according to  $N\sigma \equiv EW/-\text{error}(EW)$ , where it is important to recognize that in this case  $-\text{error}(EW)$  refers to the  $1\sigma$  (68 per cent) negative errors on an absorption lines  $EW$ , which were measured independently for each absorption line, and *not* to the 90 per cent errors ( $\approx 1.6\sigma$ ) reported in Table 2. By using the  $1\sigma$  negative errors the standard Gaussian probabilities (i.e.  $2\sigma = 95.5$  per cent,  $3\sigma = 99.73$  per cent,  $4\sigma = 99.994$  per cent) then correspond to the statistical significance of a line relative to it having an intrinsic  $EW$  of 0 eV, and thus not being detected. Fig. 11 shows that the detected lines typically have significances ranging from  $\sim 2\sigma$  through to  $\sim 10\sigma$  over a wide range in count rate; with a significance of  $\sim 30\sigma$  the extremely strong Fe xxv He $\alpha$  line in NGC 1365 (OBSID 702047010) is by far the most significant line in the sample and is omitted from the main plot for clarity.

In such a plot we would expect to see a correlation between total Fe K counts and the significance of an absorption line, where observations with higher counts in the Fe K band, and therefore better photon statistics, have stronger absorption line detections. Indeed, while there is noticeable scatter in  $N_{\text{counts}}$  versus  $\sigma$  – as would be expected given that there is a wide dispersion in absorption properties (i.e.  $N_{\text{H}}$ ,  $\xi$ ,  $v_{\text{out}}$ ) for a given total counts – a simple Spearman’s rank association shows that the data are positively correlated (Spearman’s rank coefficient,  $r_s = 0.4406$ ). This is sufficient to rule out the possibility of the two parameters being completely independent at the  $>99$  per cent significance for 34 matched pairs. Therefore,  $\sigma$



**Figure 11.** Plot of line significance,  $\sigma$  ( $\equiv EW/-\text{error}(EW)$ ), versus total counts between 5 and 10 keV. The blue circles, red squares and green crosses show the positions of the Fe xxv He $\alpha$ , Fe xxvi Ly $\alpha$  and blended Fe xxv–xxvi lines, respectively. Only absorption lines detected at  $P_{\text{MC}} \geq 95$  per cent are shown. The dashed grey band shows the  $2-3\sigma$  detection regime.

values tend to increase in higher counts observations which causes the data to gradually diverge from the  $2-3\sigma$  significance region towards the upper-right quadrant of the plot. Moreover, the vast majority of the line detections are located well away from the  $2-3\sigma$  ‘noise’ level, as would be expected if the majority of the line detections are not purely due to photon statistics.

## 7 SUMMARY AND CONCLUSIONS

Making use of data from the *Suzaku* Data Archives and Transmission System (DARTS) we have constructed broad-band spectral models for 51 Type 1.0–1.9 AGN to search for the presence of the Fe xxv He $\alpha$  and/or Fe xxvi Ly $\alpha$  absorption lines at  $E \gtrsim 6.7$  keV, robustly assessed the statistical significance of any detected absorption systems using detailed Monte Carlo simulations and probed the properties of the absorbing material using the *xstar* (v. 2.2.1bc) photoionization code. The primary findings of this work are as follows.

(i) We find that 20/51 sources (in 28/73 fitted spectra) have statistically significant Fe xxv He $\alpha$  or Fe xxvi Ly $\alpha$  absorption lines at  $E \geq 6.7$  keV in the source rest-frame, which corresponds to  $\approx 40$  per cent of the sample. 18 of the 20 Fe K absorption complexes are robustly detected at  $P_{\text{MC}} \geq 99$  per cent from Monte Carlo simulations, with those remaining only narrowly falling short of the 99 per cent criterion required for a robust detection. All absorption lines are detected independently and simultaneously in two (or more) of the available XIS detectors which further enforces the robustness of the line detections as real features.

(ii) The detected outflows fall into a range of phenomenological categories. Absorption due to a pair of Fe xxv He $\alpha$  and Fe xxvi Ly $\alpha$  from the same photoionized absorber account for 7/20 of the detected outflows (in 12/28 spectra) making them the most frequently observed form of highly ionized outflow in the sample. Absorption due to solely Fe xxv He $\alpha$  or Fe xxvi Ly $\alpha$  is less frequently

observed, being detected in two sources (3 spectra) and nine sources (10 spectra), respectively. The remaining two sources have outflows which are best fitted by two unresolved and blended Fe XXV–XXVI absorption systems with different outflow velocities.

(iii) By fitting the absorption with the *xstar* photoionization code we find that the absorbers are characterized by  $21.5 < \log(N_{\text{H}}/\text{cm}^{-2}) \leq 24$  and  $2.5 < \log(\xi/\text{erg cm s}^{-1}) \leq 6$ , with mean values of  $\log(N_{\text{H}}/\text{cm}^{-2}) \approx 23$  and  $\log(\xi/\text{erg cm s}^{-1}) \approx 4.5$ , respectively. The distribution of outflow velocities covers a wide range, from  $v_{\text{out}} < 1500 \text{ km s}^{-1}$  up to  $100\,000 \text{ km s}^{-1}$ , with 90 per cent of the observed absorbers having  $v_{\text{out}} > 1500 \text{ km s}^{-1}$  which makes them systematically faster than typical absorbers found in the soft X-ray band. The median outflow velocity is of the order of  $\sim 0.056 c$ . Moreover, 60 per cent of the absorbers also have  $v_{\text{out}} > 10\,000 \text{ km s}^{-1}$  and are thus consistent with the so-called UFOs defined by T10A. A K-S test shows that the overall distributions for all three parameters are not statistically distinguishable from the analogous distributions measured by T10A and T11 using *XMM–Newton*.

Overall, the results of this work are consistent with those obtained using *XMM–Newton*, and, combined, the two studies provide strong evidence for the presence of very highly ionized, often high-velocity, outflows in the central region of a large fraction of not only radio-quiet AGN, but also suggests that they may be prevalent in radio-loud sources as well. The possible prevalence of very highly ionized, high-velocity winds in AGN is consistent with theoretical models which argue that such outflows are an important phenomenon which may play a role in Galactic-scale feedback scenarios. We return to this point in a forthcoming work (Paper II; Gofford et al., in preparation) where we probe the global energetics of the detected absorbers, assess for any correlations between the outflow parameters and those of the central SMBH, and discuss both the origins of the absorbing material and its likely launching mechanism.

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## SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article:

**Appendix A.** Summary of observations.

**Appendix B.** Ratio and contour plots.

**Appendix C.** Modelling Complexities.

**Appendix D.** Full model details

(<http://mnras.oxfordjournals.org/lookup/suppl/doi:10.1093/mnras/sts481/-/DC1>).

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